

# Stellar and Planetary Parameters for K2's Late-type Dwarf Systems from C1 to C5

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### Abstract

The NASA K2 mission uses photometry to find planets transiting stars of various types. M dwarfs are of high interest since they host more short-period planets than any other type of main-sequence star and transiting planets around M dwarfs have deeper transits compared to other main-sequence stars. In this paper, we present stellar parameters from K and M dwarfs hosting transiting planet candidates discovered by our team. Using the SOFI spectrograph on the European Southern Observatory's New Technology Telescope, we obtained  $R \approx 1000 J$ -, H-, and K-band (0.95–2.52 µm) spectra of 34 late-type K2 planet and candidate planet host systems and 12 bright K4– M5 dwarfs with interferometrically measured radii and effective temperatures. Out of our 34 late-type K2 targets, we identify 27 of these stars as M dwarfs. We measure equivalent widths of spectral features, derive calibration relations using stars with interferometric measurements, and estimate stellar radii, effective temperatures, masses, and luminosities for the K2 planet hosts. Our calibrations provide radii and temperatures with median uncertainties of 0.059  $R_{\odot}$  (16.09%) and 160 K (4.33%), respectively. We then reassess the radii and equilibrium temperatures of known and candidate planets based on our spectroscopically derived stellar parameters. Since a planet's radius and equilibrium temperature depend on the parameters of its host star, our study provides more precise planetary parameters for planets and candidates orbiting late-type stars observed with K2. We find a median planet radius and an equilibrium temperature of approximately 3  $R_{\oplus}$  and 500 K, respectively, with several systems (K2-18b and K2-72e) receiving near-Earth-like levels of incident irradiation.

Key words: methods: data analysis – planetary systems – stars: fundamental parameters – stars: late-type – techniques: spectroscopic

#### 1. Introduction

Small, low-luminosity M dwarfs are the most common type of star in the Galaxy, but their properties are less well understood than those of hotter solar-type stars. There are still significant discrepancies between theoretical models and observations of M dwarf spectra (e.g., Hoeijmakers et al. 2015), and we are still uncertain as to why the occurrence rate of small, short-period planets is higher for M dwarfs and the occurrence rate of gas giants (on both close and wide orbits) is lower for M dwarfs when compared to solar-like stars, as shown in studies of the Kepler field (Dressing & Charbonneau 2013; Gaidos et al. 2014; Morton & Swift 2014; Dressing & Charbonneau 2015; Muirhead et al. 2015) and other surveys (Shields et al. 2016). There are a few exceptions to the low occurrence rate of gas giants around M dwarfs; there has been at least one confirmed gas giant orbiting an M dwarf (Johnson et al. 2012).

Fortunately, the discovery of exoplanets around M dwarfs is much easier when compared to finding exoplanets around Sunlike stars. For example, while a transiting  $2R_{\oplus}$  planet would have a transit depth of 0.03% when orbiting the Sun, that same planet would have a transit depth of 0.5% for an M5 dwarf. Using planet candidates from the original Kepler mission, Howard et al. (2012) and Mulders et al. (2015a, 2015b) showed that the occurrence rates of small planets are higher for M dwarfs than for any other type of main-sequence star. Other surveys, such as MEarth (Charbonneau et al. 2009; Berta-Thompson et al. 2015) and Transiting Planets and Planetesimals Small Telescope, have also successfully identified interesting new planets transiting M dwarfs (Gillon et al. 2016). Additionally, M dwarfs provide our best chances to identify nearby potentially habitable planets, since the

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habitable zone around M dwarfs, when compared to those around other main-sequence stars, is closer to the M dwarf, due to its lower luminosity. This is exemplified by the discovery of Proxima Centauri b, a small, likely temperate planet orbiting the closest star to the Sun (Anglada-Escudé et al. 2016; Damasso & Del Sordo 2016).

Host star properties must be well understood in order to be able to derive planet properties. Unfortunately, the stellar properties of M dwarfs are challenging to predict from photometry (due to M dwarfs being intrinsically faint and the modeling uncertainties as described above and by Mann et al. 2015). The most accurate parameters of M dwarfs are derived from interferometric data (Boyajian et al. 2012b) or photometric and spectroscopic observations of double-lined eclipsing binaries (Torres et al. 2010).

For systems where such observations are not feasible, several authors have developed a calibration method based on medium-resolution, near-infrared spectra in order to infer the stellar properties of these M dwarfs from empirical observations (Mann et al. 2015; Newton et al. 2015; Terrien et al. 2015) and stellar models (Rojas-Ayala et al. 2012), while others have applied similar empirical calibration techniques to the optical part of the spectrum (Neves et al. 2014; Maldonado et al. 2015). By measuring the equivalent widths (EWs), or the strength of any given absorption feature, one can calculate stellar parameters by calibrating from a reference sample with previously measured parameters of interest. Since the EW of an absorption feature varies with photospheric temperature and surface gravity, this approach allows these parameters (and related quantities, like stellar radius and mass) to be calculated.

Using the repurposed *Kepler* spacecraft, the *K*2 mission is continuing to observe many stars in the Galaxy in the search for more exoplanets (Howell et al. 2014). However, *K*2 has some limitations. With just two (out of four) operating reaction wheels, the spacecraft can observe only along the ecliptic plane with observation windows of 80 days per campaign. None-theless, *K*2 has provided astronomers with powerful data enabling a large number of candidate and confirmed exoplanets (Vanderburg & Johnson 2014; Crossfield et al. 2015; 2016; Foreman-Mackey et al. 2015; Huang et al. 2015; Montet et al. 2015; Sanchis-Ojeda et al. 2015; Sinukoff et al. 2015).

In this paper we analyze medium-resolution, near-infrared spectra of candidate planetary systems detected by K2 to provide updated stellar and planetary parameters. We measure EWs to infer stellar radii and effective temperatures, and subsequently planetary radii and equilibrium temperatures. In Section 2 we briefly explain our target selections and how we compiled our planet candidate list. In Section 3 we describe our observational techniques, data reduction, and various calibration samples. In Section 4 we explain the process by which we obtain our stellar and planetary parameters and compare our derived stellar parameters with those of previously spectroscopically and interferometrically measured stellar parameters. In Section 5 we summarize our results and describe future work relevant to this paper.

## 2. Target Selection and Planet Candidate Search

We initially selected our K2 M dwarf candidates from Campaigns 1 through 5. Our team selected and proposed latetype dwarf targets to the K2 mission as described by Crossfield et al. (2016). In brief, we selected targets as being likely lowmass dwarfs by a combined color and proper motion cut with (V-J) > 2.5,  $V + 5 \log \mu + 5 < 10$ , and  $(6V - 7J - 3) < 5 \log \mu$  (where  $\mu$  is the proper motion; Crossfield et al. 2015). The combination of the color and proper motion cut greatly reduces giants from our sample and further narrows down the M dwarf candidate list. Finally, we imposed a magnitude limit of Kp < 16.5 mag (Crossfield et al. 2016).

We further identify likely low-mass planet-hosting dwarf stars, as explained in Crossfield et al. (2016). In brief, we used the TERRA algorithm (Petigura et al. 2013) to search for planet transits that have a signal-to-noise ratio (S/N) > 12, which are called threshold-crossing events (TCEs). TCEs are required to have orbital periods of  $P \ge 1$  day and to have at least three transits. These restrictions, along with the diagnostic tests that TERRA provides, show whether the object is a candidate transiting planet, binary star system, another variable object, or noise. If a planet candidate is found, TERRA is iteratively repeated after removing the identified transit signals (described by Sinukoff et al. 2016) to see whether there are any additional planets in the system.

# 3. Observations

We acquired our infrared spectra at the 3.58 m European Southern Observatory (ESO) New Technology Telescope (NTT) using the SOFI spectrograph (Moorwood et al. 1998) as part of program 194.C-0443 (PI: I. J. M. Crossfield). We observed through 13 full or partial usable nights in 2015 and 2016. We used two grisms, red and blue, to produce a total spectrum for each object spanning a continuous wavelength<sup>20</sup> range from 0.95 to 2.52  $\mu$ m at a resolution of  $R \approx 1000$ . Dome flats and lamps were taken either at the start or at the end of each observing night. Our observation sample comprises 34 stars observed by *K2* in fields 1 through 5, along with 12 bright K and M dwarfs with interferometrically measured stellar parameters (refer to Table 1 for our calibration sample).

For all observations, we used an ABBA nodding pattern to obtain the spectrum of the object, while removing the spectrum of the background, including sky emission lines and dark current. The exposure times for each frame range from the minimum allowed exposure time (1.182 s) to 120 s. We typically took at least six separate spectra (for each grism) for all the targets. Either immediately before or after each M dwarf candidate, we observed a nearby A0V star for telluric corrections. If the observation for one grism took more than 10 minutes, its A0V calibrator would be taken before the start of the first grism and then taken again after the second grism exposure had finished, for their respective grisms. We identified suitable A0V stars using the IRTF's online tool.<sup>21</sup>

## 3.1. Data Reduction

The raw data taken at the NTT were reduced by using a combination of Python, Image Reduction and Analysis Facility (IRAF) software,<sup>22</sup> and using various Interactive Data Language (IDL) programs. We flat-fielded the raw spectra in order to correct for any pixel-to-pixel variation. Wavelength calibrations were done by taking Xe arc spectrum for both

<sup>&</sup>lt;sup>20</sup> The blue grism spans the wavelength range from 0.95 to 1.64  $\mu$ m, while the red grism spans the wavelength range from 1.53 to 2.52  $\mu$ m. Note that there is a small overlap from both grisms in the *H* band, thus allowing the fully reduced spectra of all of our stars to be continuous.

<sup>&</sup>lt;sup>21</sup> http://irtfweb.ifa.hawaii.edu/cgi-bin/spex/find\_a0v.cgi

<sup>&</sup>lt;sup>22</sup> Developed at the National Optical Astronomy Observatory.

	Table 1	
Stellar	Calibration	Sample

Star	SpT <sup>a</sup>	$R_*$	$T_{\rm eff}$	$L_*$	Reference	Notes
		$(R_{\odot})$	(K)	$(L_{\odot})$		
GJ 176	M2.5V	0.453(22)	3679(77)	0.0337(43)	von Braun et al. (2014)	
GJ 205	M1.5V	0.5735(44)	3801(9)	0.0616(11)	Boyajian et al. (2012b)	
GJ 436	M3V	0.455(18)	3416(53)	0.0253(25)	von Braun et al. (2012)	1
GJ 526	M1.5V	0.4840(84)	3618(31)	0.0360(18)	Boyajian et al. (2012b)	
GJ 551	M5.5V	0.1410(70)	3054(79)	0.00155(22)	Boyajian et al. (2012b)	
GJ 570A	K4V	0.739(19)	4507(58)	0.202(15)	Demory et al. (2009)	
GJ 581	M2.5V	0.299(10)	3442(54)	0.0113(10)	von Braun et al. (2011)	2
GJ 699	M4.0V	0.1869(12)	3222(10)	0.003380(60)	Boyajian et al. (2012b)	
GJ 702B	K5Ve	0.6697(89)	4400(150)	0.150(46)	Boyajian et al. (2012b)	
GJ 845	K5V	0.7320(60)	4555(24)	0.207(34)	Demory et al. (2009)	
GJ 876	M3.5V	0.3761(59)	3129(19)	0.0122(39)	von Braun et al. (2014)	3
GJ 880	M1.5V	0.5477(48)	3713(11)	0.0512(90)	Boyajian et al. (2012b)	

Note.

<sup>a</sup> Spectral types were adopted from the interferometric works, with the following exceptions: (1) Kirkpatrick et al. (1991); Hawley et al. (1996); (2) Henry et al. (1994); and (3) were linearly interpolated from Pickles (1998).

grisms either at the beginning or at the end of the night. Using IRAF, emission lines from the taken Xe arc frames were manually selected by comparing them to the SOFI manual.<sup>23</sup> One-dimensional spectra were then extracted for identifying the star's spectrum. IRAF had difficulty tracing the 2D spectra of our fainter targets, so for these stars we used brighter stars during that night to define a static extraction aperture.

We subsequently used the IDL routines of Vacca et al. (2003) to process our spectra. First, with xcombspec (from the SpeXtool software package by Cushing et al. 2004), we combined multiple exposures for a given grism of an object into one spectrum. Any spectra that are not shown to have similar spectral features with the other exposures for that star and grism were excluded.

We corrected for telluric absorption by using our AOV spectra with the xtellcor\_general routine. Spectra of A0V stars were used since these stars are mostly composed of featureless spectra, with the exception of hydrogen absorption. Differences between the hydrogen lines in the A0V and a model Vega spectrum were corrected for, and then the object's spectrum was divided by the resulting telluric spectrum of the A0V; the observations for the telluric calibrator were usually taken within a short time (approximately 15 minutes) and have a similar airmass (within 0.3 airmass) to the object (Rojas-Ayala et al. 2012). We note that for some of the observations, the telluric calibrator's spectrum was sufficiently different from that of Vega that some residual H lines remain in the M dwarf candidate's spectrum. Additionally, the large differences in airmass left residual telluric features in some of the spectra, and any spectra that were contaminated were removed from our analysis.

The last step for the reduction process was to combine the two different grisms using xmergexd. We then used several strong absorption features in each spectrum to correct for radial velocity (RV) shifts and/or offsets in our wavelength calibration. Finally, we interpolated all spectra to put them on the same wavelength scale. All of the objects in our sample have a S/N that ranged from 20 (for the faint K2 targets)<sup>24</sup> to

over 200 (for the brighter, interferometric calibration targets). We show a representative reduced spectrum in Figure 1.

# 3.2. Calibration Sample

We applied the relations from a variety of works, such as Neves et al. (2014) and Maldonado et al. (2015), which fit various functions for a variety of EW ratios, and Terrien et al. (2015), which measured *H*-band atomic features, to stars with previously measured radii and/or effective temperatures. However, stars that are interferometrically measured are preferred to these samples since measurements from interferometry are more accurate and precise when compared to spectroscopic, EW-based methods. Although most interferometrically measured stars lie too far north to be observed with SOFI, we managed to obtain spectra of 12 stars with previously interferometrically determined stellar radii and effective temperatures. These stars form our calibration sample, and their properties are summarized in Table 1.

#### 4. Spectral Analysis

Mould (1976) was the first to use infrared absorption line strengths to estimate the radii and effective temperatures of low-mass dwarfs. The strengths of absorption features corresponding to a given element or molecule depend on the effective temperatures of the star. Changing the temperature of the star then changes the electronic (or vibrational) population levels of the element (or molecule) in the M dwarf atmosphere. M dwarf radii are related to their effective temperatures so that they roughly follow a linear relation from 4700 K and 0.7  $R_{\odot}$  down to at least 3300 K and 0.3  $R_{\odot}$ . Some of the absorption features in the spectrum can also present information about the stellar surface gravity. The lines of alkali elements, for example, are affected by surface gravity and can then be used to distinguish old dwarf stars, young dwarf stars, and giants with similar temperatures (Spinrad 1962; Steele & Jameson 1995; Lyo et al. 2004; Schlieder et al. 2012).

The EW is defined by the following equation:

$$\mathrm{EW}_{\lambda} = \int_{\lambda_{\mathrm{I}}}^{\lambda_{\mathrm{2}}} \left[ 1 - \frac{F(\lambda)}{F_{c}(\lambda)} \right] d\lambda, \tag{1}$$

<sup>&</sup>lt;sup>23</sup> Provided by ESO.

 $<sup>^{24}</sup>$  K2 targets that had a S/N of 20 were removed from the likely low-mass dwarf list, thus making our final 34-star sample.



**Figure 1.** Sample spectrum of one of our *K2* targets (EPIC 201367065 or K2-3) that covers a continuous wavelength from 0.95 to 2.52  $\mu$ m and is normalized to the median flux value. Note that we ignore regions heavily contaminated by telluric features (e.g., wavelength ranges that are within 1.35–1.45  $\mu$ m and 1.80–1.95  $\mu$ m). After data reduction is complete, we trim an approximate 0.01–0.02  $\mu$ m off the edges of the wavelength ranges. Spectra of all our stars are available as electronic supplements to this paper.

where  $F(\lambda)$  is the flux of the absorption feature between  $\lambda_1$  and  $\lambda_2$ , and  $F_c(\lambda)$  is the continuum flux. We investigate the features used by Cushing et al. (2005), Rojas-Ayala et al. (2012), Newton et al. (2014), and Newton et al. (2015) for our work. The features, shown in Table 2, are slightly adjusted owing to differences in resolution of the spectrographs—typically our integration ranges are slightly wider than those previously presented. Additionally, any spectral line doublets and molecular bands used in our empirical indices are treated as single features in the EW calculations. The blue continuum and red continuum of each feature are also adjusted such that they would not overlap with any nearby feature windows. In the following sections, we describe the steps that are taken to infer the stellar and planetary parameters using these EW measurements of our *K2* and calibration samples.

## 4.1. Spectral Classification

We visually estimated the spectral types (SpT) of each of our stars by comparing our SOFI spectra to spectra of standard stars in the IRTF Spectral Library (Cushing et al. 2005; Rayner et al. 2009). Then, we convolved the library spectra from G8V to M7V down to the resolution of SOFI and plotted these against each of our SOFI spectra. We estimated each SpT and a corresponding uncertainty three times by independently comparing spectra in the *J*-, *H*-, and *K*- bandpasses. The final uncertainty on each SpT corresponds to the uncertainty on the weighted mean and thus represents our best estimate of the error on this quantity. We then compute a single SpT for each star using a weighted mean. The SpT and uncertainty are rounded to the nearest 1/10 of a type. Out of the 34 stars in our *K2* sample, we identify 27 as M dwarfs.

During our visual spectral inspection, we compared our spectra to the library spectra of giant stars in order to remove giants as early as possible in our analysis process. We identified only one star as a likely giant: EPIC 202710713, which Huber

et al. (2016) and Dressing et al. (2017) also classified as an evolved star.

#### 4.2. Stellar Parameters

For each absorption feature and stellar parameter (radius and effective temperature), we use least-squares fitting to determine the dependence of those parameters on the EWs calculated from the spectra. Various functional forms of EWs are used to fit the calibration sample's parameters. They include all combinations of linear, quadratic, and a ratio of EWs of two different absorption features. For example, in the simplest linear case, one lets the EW for the chosen absorption feature be the independent variable, while stellar radius or effective temperature is the dependent variable. After calculating the linear term and the offset, one then uses all the EWs to calculate the stellar radius for all the stars in our sample. This process is then repeated for all the absorption features in the spectra, all the stellar parameters, each calibration sample, and each functional combination of EWs. To account for intrinsic scatter in stellar properties, we include an additional noise term, tuned to give  $\chi^2_{\rm red} \approx 1$  in the best cases. We find that scatter terms of 100 K and  $0.05 R_{\odot}$  fulfill this criterion.

In order to find the optimal fit for each calibration sample, we then select the model giving the lowest Bayesian information criterion (BIC) value and the lowest scatter in the fit residuals. We use a Monte Carlo approach to estimate the uncertainties on the fit coefficients and inferred stellar parameters. Random Gaussian distributions are then used to generate synthetic data sets of EWs, stellar radii, and effective temperatures. A total of 1000 trials are used for calculating the uncertainties for each parameter.

## 4.3. Calibration Relations and Literature Comparison

Because some of our spectra contain residual systematics near prominent H lines, we find only poor fits using EWs located near these lines (Brackett 11-21). Viewing all possible combinations of the remaining EWs, we determine that the optimal fits for calculating our parameters are determined by having a low BIC value for the fit and comparing it to the median uncertainty of all the uncertainties in a given combination of EWs. We present the following equations for calculating stellar radius and effective temperature:

$$\frac{T_{\rm eff}}{K} = a + b \left(\frac{{\rm Mg}_{\rm 1.57}}{{\rm Al}_{\rm 1.31}}\right) + c \left(\frac{{\rm Al}_{\rm 1.67}}{{\rm Ca}\,{\rm I}_{\rm 1.03}}\right),\tag{2}$$

$$\frac{R_*}{R_{\odot}} = a + b \,(\mathrm{Mg}_{1.57}) + c \left(\frac{\mathrm{CO}_{2.29}}{\mathrm{Na}\,\mathrm{I}_{1.14}}\right). \tag{3}$$

Table 3 lists the best-fitting coefficients and the covariance matrix for each fit. Note that some coefficients exhibit significant correlations, suggesting that uncertainties would be underestimated if these correlations were neglected.

Based on the range of our calibration sample, we restrict ourselves to stars in the range  $3000 \text{ K} < T_{\text{eff}} < 4500 \text{ K}$  and  $0.2 < R_*/R_{\odot} < 0.7$ . There is overall excellent agreement between our derived values for radius and effective temperature, while four stars (GJ 551, GJ 699, GJ 526, GJ 876) have somewhat larger deviations in stellar radius and/or effective temperature. Figures 2 and 3 compare the inferred and literature values for our calibrated sample. The middle and bottom panels of these two figures show that the dispersions of the residuals

Table 2	
J-, H-, and K-band Equivalent Width	Features

Feature	Feature	Window	Blue Co	ontinuum	Red Co	Red Continuum		
- Culture	(μ	m)	(μ	em)	(µm)			
Ca I (1.03 μm)	1.0320	1.0365	1.0280	1.0315	1.0368	1.0377		
Na I (1.14 μm)	1.1361	1.1432	1.1270	1.1327	1.1478	1.1572		
Al (1.31 µm)	1.3125	1.3180	1.3060	1.3090	1.3180	1.3220		
Mg (1.48 μm)	1.4865	1.4905	1.4810	1.4850	1.4920	1.4960		
Mg (1.50 μm)	1.5002	1.5075	1.4910	1.4983	1.5090	1.5163		
Mg (1.57 μm)	1.5725	1.5797	1.5665	1.5720	1.5810	1.5865		
Si (1.58 µm)	1.5875	1.5925	1.5820	1.5865	1.5930	1.5975		
CO (1.62 µm)	1.6178	1.6280	1.6048	1.6150	1.6300	1.6402		
Al (1.67 µm)	1.6698	1.6790	1.6558	1.6650	1.6800	1.6892		
Mg (1.71 μm)	1.7089	1.7139	1.7000	1.7050	1.7149	1.7199		
Na I (2.20 μm)	2.2020	2.2120	2.1890	2.1990	2.2125	2.2225		
Ca I (2.26 µm)	2.2586	2.2696	2.2480	2.2570	2.2700	2.2800		
CO (2.29 μm)	2.292	2.315	2.286	2.290	2.315	2.320		

Note. All the wavelengths are presented at their rest wavelength.

are 0.059  $R_{\odot}$  (16.09%) and 160 K (4.33%) for stellar radius and effective temperature, respectively. All of the stars in our calibration sample, with the exception of GJ 526, have published luminosities (calculated using the Stefan–Boltzmann law) within 1 $\sigma$  of our inferred values. Finally, we estimate each star's mass by inverting the mass–radius relationship of Maldonado et al. (2015). The full set of stellar values is listed in Table 4, and the *K2* stellar parameters are plotted in Figure 4.

We also independently compare our stellar parameters to those of Dressing et al. (2017). Out of our 34-star K2 sample (as referenced in Table 4), we share 21 stars in common with their sample. While this work calculates stellar parameters using the spectra acquired with NTT/SOFI, Dressing et al. (2017) use two different instruments in their work. The SpeX instrument, on the NASA Infrared Telescope Facility, provides wavelength coverage from 0.7 to 2.55  $\mu$ m at a resolution of  $R \approx 2000$  (Rayner et al. 2003). The other instrument used was TripleSpec on the Palomar 200<sup>"</sup>, providing wavelength coverage from 1.0 to 2.4  $\mu$ m at a resolution of  $R \approx 2500-2700$ (Herter et al. 2008). Dressing et al. (2017) derive and compare stellar parameters using EW-based relations developed by Newton et al. (2015) and index-based relations from Mann et al. (2013). Both sets of relations were calibrated using a set of stars with interferometrically determined parameters from Boyajian et al. (2012b). Ultimately, effective temperatures, stellar radii, and luminosities were derived using the Newton et al. (2015) relations, stellar masses<sup>25</sup> and metallicity were calculated using the Mann et al. (2013) relations, and surface gravities were calculated from masses and stellar radii.

Comparing the parameters derived by Dressing et al. (2017) with those shown in Figures 5 and 6, we find  $\chi^2_{red} < 1$  in both cases. This indicates that there is an excellent agreement between our two methods and verifies the validity of our approach. Additionally, our stellar parameters are consistent with those from a number of previous publications (Crossfield et al. 2015; Montet et al. 2015; Petigura et al. 2015; Mann et al. 2016; Obermeier et al. 2016; Schlieder et al. 2016).

The most highly discrepant system evident in Figure 6 seems to be the effective temperature of EPIC 211770795. Our estimate is significantly lower than the 4750 K estimated by Dressing

et al. (2017). Their value is larger than the 4500 K upper limit determined from our calibration sample (see Figure 2), providing further evidence that our relations are not well calibrated beyond this range. Furthermore, we see an offset between our effective temperature values and those reported by Dressing et al. (2017), demonstrating that systematic calibration errors may still play a role in one or both of these analyses. As seen with the indexbased relations of Mann et al. (2015), our EW-based relations also start to saturate around 4000 K and could systematically effect any derived planetary parameters, such as equilibrium temperatures.

Metallicity could be a factor for some stars and could cause a shift in effective temperature and stellar radius. The larger uncertainties in our stellar parameters when compared to those of Dressing et al. (2017) may result from a range of stellar metallicities.

Additionally, we compare all 34 of our stellar parameters with the photometrically derived stellar parameters from Huber et al. (2016), shown in Figures 7 and 8. Figure 7 shows that there is a median increase of  $0.15 R_{\odot}$  when comparing our stellar radii to those of Huber et al. (2016). Figure 8 shows a general agreement in effective temperature between both of our works, with the exception of EPIC 204489514 and EPIC 205145448.

We note that the analysis done in Huber et al. (2016) is subject to the limitations of broadband photometry. Furthermore, Huber et al. (2016) note that model-based estimates tend to underpredict stellar radii by 20% (Boyajian et al. 2012a) and encourage the use of empirical calibrations for estimating the stellar parameters in cool dwarfs. Lastly, our empirically calculated parameters are in agreement with those in Dressing et al. (2017) for the points where we disagree with the values of Huber et al. (2016), giving us further confidence in our results.

#### 4.4. Planetary Parameters

Radii and equilibrium temperatures of transiting planets are calculated using the stellar parameters of its host star. Using the transit depths and periods measured using K2 photometry (Crossfield et al. 2016) and our newly calculated stellar parameters, planet radii are determined with the following

<sup>25</sup> Using the effective temperatures from the Newton et al. (2015) relation.

Quantity	Formula	a	b	С
T <sub>eff</sub>	$a + b (Mg_{1.57}/Al_{1.31}) + c (Al_{1.67}/Ca I_{1.03})$	2989.5	-577.05	53.804
	Uncertainties:	78.56147	52.42034	10.44419
	Covariance:			
		6171.9	-3355.2	-493.87
		-3355.2	2747.9	265.68
		-493.87	265.68	109.08
<i>R</i> *	$a + b (Mg_{1.57} \text{ Å}) + c (CO_{2.29}/Na I_{1.14})$	0.18552	1265.2	0.010852
	Uncertainties:	0.02569482	117.2119	0.005063553
	Covariance:			
		0.00066022	-1.8673	-0.00003695
		-1.8673	13739	-0.2305
		-0.00003695	-0.2305	0.00002564

Table 3Equivalent Width Formulae



**Figure 2.** Stellar radius for the stars in our interferometric calibration sample, from the literature (blue circles) and derived using Equation (3) (red circles). The middle and bottom panels show the absolute and fractional deviations for each star. The dispersion of the residuals is  $0.059 R_{\odot}$  and 16.09%, respectively. Our sample spans from 0.2 to  $0.7 R_{\odot}$ .



**Figure 3.** Stellar effective temperature for the stars in our interferometric calibration sample, from the literature (blue circles) and derived using Equation (2) (red circles). The middle and bottom panels show the absolute and fractional deviations for each star. The dispersion of the residuals is 160 K and 4.33%, respectively. Our sample spans from 3000 to 4500 K.

equation:

$$\Delta L = \left(\frac{R_p}{R_*}\right)^2,\tag{4}$$

where  $\Delta L$  is the transit depth of the planet candidate with respect to its host star.

Calculating the equilibrium temperature of a planet candidate requires more parameters from the planet and its host star. The following equation calculates the equilibrium temperatures for each *K*2 planet or candidate as a comparison to our own Earth–Sun system:

$$T_{\rm eq} = (270K) \left( \frac{T_{\rm eff}}{T_{\rm eff,\odot}} \right) \left[ \left( \frac{R_*}{R_\odot} \right) \left( \frac{1 \text{ au}}{a} \right) \right]^{1/2}, \tag{5}$$

where  $T_{\rm eff}$  is the effective temperature of the star,  $R_*$  is the radius of the star, and *a* is the semimajor axis of the planet orbiting its parent star. Here we calculate the semimajor axis of the planet by using Kepler's third law. The 270 K equilibrium temperature scaling factor corresponds to a Bond albedo of 0.3, which is comparable to that inferred for gas giants more highly irritated than Earth. All uncertainties are propagated through the entire calculation for planet radii. We present the derived values for our *K2* planets and planet candidates in Table 5 and plot these derived values (along with incident irradiation) in Figure 9.

Our sample shown in Figure 9 includes 18 validated planets and 19 remaining planet candidates. While fitting for the lightcurve parameters of these remaining candidates, degeneracies (such as impact parameters near unity) arose that preclude any precise determination of  $R_p/R_*$ . The candidates have much larger uncertainties on their size, which typically makes statistical validation much more difficult. Based on the paucity of large (>6  $R_{\oplus}$ ) planets orbiting M dwarfs (Johnson et al. 2007, 2010), the  $\gtrsim$ 9 candidates larger than this size are likely false positives; since planet validation is not the aim of this work, we retain the previously assigned designation of planet candidate.

In addition to these likely false positives, our validated planets include several hot Neptunes and two planets (K2-18b and K2-72e) that lie near the habitable zone. Of our whole K2 sample, only eight planets (three of which are still planet candidates) are smaller than 1.6 Earth radii. According to Rogers (2015), planets smaller than 1.6 Earth radii are likely to

Table 4       Derived Stellar Parameters												
Star <sup>a</sup>	SpT	$egin{array}{c} R_{*} \ (R_{\odot}) \end{array}$	T <sub>eff</sub> (K)	$M_*$ $(M_{\odot})$	$L_*$ $(L_{\odot})$	$\alpha_{2000}$	$\delta_{2000}$	Kp (mag)	J (mag)	H (mag)	K (mag)	Notes <sup>b</sup>
GJ 176	M2.4(0.7)	0.528(53)	3605(170)	0.546(55)	0.0423(63)	04 <sup>h</sup> 42 <sup>m</sup> 55 <sup>§</sup> 774	$+18^{\circ}57'29''_{\cdot}404$		6.462	5.824	5.607	1
GJ 205	M1.4(1.1)	0.645(56)	3735(172)	0.665(55)	0.0726(96)	05 <sup>h</sup> 31 <sup>m</sup> 27 <sup>s</sup> 395	-03°40'38."031		4.83	4.05	3.90	2
GJ 436	M2.5(0.6)	0.469(53)	3630(165)	0.484(57)	0.0343(57)	11 <sup>h</sup> 42 <sup>m</sup> 11 <sup>s</sup> 094	+26°42′23.″65		6.900	6.319	6.073	3
GJ 526	M2.3(1.0)	0.457(53)	3647(173)	0.472(57)	0.0332(56)	13 <sup>h</sup> 45 <sup>m</sup> 43 <sup>s</sup> .776	+14°53′29.″463		5.18	4.78	4.415	2
GJ 551	M4.8(0.6)	0.150(59)	2887(234)	0.104(78)	0.00141(79)	14 <sup>h</sup> 29 <sup>m</sup> 42 <sup>s</sup> .948	-62°40′46″163		5.357	4.835	4.384	4
GJ 570A	K4.1(0.7)	0.615(55)	4498(417)	0.636(55)	0.139(22)	14h57m28 s 001	-21°24′55″713		3.83	3.23	3.10	4
GJ 581	M2.5(0.7)	0.279(60)	3624(208)	0.266(72)	0.0121(37)	15 <sup>h</sup> 19 <sup>m</sup> 26 <sup>s</sup> .823	-07°43′20″21		6.706	6.095	5.837	5
GJ 699	M3.1(0.8)	0.268(55)	3182(185)	0.253(68)	0.0066(20)	17 <sup>h</sup> 57 <sup>m</sup> 48 <sup>s</sup> .498	+04°41′36.″207		5.244	4.83	4.524	2, 6
GJ 702B	M3.5(0.6)	0.641(56)	4009(225)	0.661(56)	0.095(13)	18h05m27 <sup>s</sup> .421	+02°29′56.″42					2
GJ 845	M1.1(2.3)	0.74(16)	4565(706)	0.76(15)	0.216(72)	22 <sup>h</sup> 03 <sup>m</sup> 21 <sup>s</sup> .658	-56°47′09."516		2.894	2.349	2.237	4
GJ 876	M2.4(0.6)	0.291(54)	3309(170)	0.282(65)	0.0091(24)	22h53m16.5733	-14°15′49″318		5.934	5.349	5.010	1
GJ 880	M0.7(1.0)	0.571(54)	3897(193)	0.591(55)	0.0676(97)	22h56 <sup>m</sup> 34 <sup>s</sup> 804	+16°33'12."354		5.360	4.800	4.523	2
201205469	M0.9(1.0)	0.559(57)	3923(198)	0.577(58)	0.066(10)	11 <sup>h</sup> 16 <sup>m</sup> 28 <sup>s</sup> .114	-03°58'31."58	14.887	12.422	11.712	11.577	
201208431	K7.7(1.2)	0.658(56)	3900(195)	0.678(55)	0.090(12)	11 <sup>h</sup> 38 <sup>m</sup> 58 <sup>s</sup> 954	-03°54′20″11	14.409	12.367	11.747	11.571	
201367065	M0.1(1.1)	0.565(61)	3976(205)	0.584(63)	0.072(12)	11 <sup>h</sup> 29 <sup>m</sup> 20 <sup>s</sup> 388	-01°27′17″23	11.574	9.421	8.805	8.561	
201465501	M2.8(0.6)	0.366(53)	3460(164)	0.369(61)	0.0173(36)	11 <sup>h</sup> 45 <sup>m</sup> 03 <sup>s</sup> 472	+00°00'19"08	14.957	12.451	11.710	11.495	
201617985	M0.5(0.7)	0.606(39)	3853(135)	0.626(39)	0.0975(70)	11 <sup>h</sup> 57 <sup>m</sup> 57 <sup>s</sup> 998	+02°19′17″31	14.110	11.719	11.094	10.900	8
201690311	K4.2(1.2)	0.697(94)	3948(203)	0.714(92)	0.106(21)	11 <sup>h</sup> 49 <sup>m</sup> 16 <sup>s</sup> 849	+03°28'32"05	15.288	13.463	12.873	12.729	
201717274	M3.0(0.8)	0.368(80)	3528(165)	0.373(92)	0.0188(59)	11h35m18 5664	$+03^{\circ}56'02''96$	14.828	12.911	12.367	12.153	
201912552	M3.0(0.9)	0.411(53)	3527(162)	0.419(58)	0.0234(44)	11 <sup>h</sup> 30 <sup>m</sup> 14 <sup>s</sup> 510	$+07^{\circ}35'18''21$	12.473	9.763	9.135	8.899	
204489514	M2.7(1.3)	0.230(56)	3096(198)	0.207(71)	0.0044(15)	16 <sup>h</sup> 03 <sup>m</sup> 01 <sup>s</sup> 616	-22°07′52″40	14.080	12.731	12.110	11.729	
205145448	M1.8(1.9)	0.402(52)	4035(259)	0.409(59)	0.0384(75)	16 <sup>h</sup> 33 <sup>m</sup> 47 <sup>s</sup> 672	-19°10′40″04	13.651	10.977	10.351	10.120	
205916793	M0.0(0.8)	0.707(60)	4103(230)	0.724(56)	0.127(17)	$22^{h}32^{m}13^{s}004$	-17°32′38″38	13.441	11.850	11.231	11.075	
205924614	K4.2(1.2)	0.769(63)	4240(259)	0.785(59)	0.172(22)	22 <sup>h</sup> 15 <sup>m</sup> 00 <sup>s</sup> 462	-17°15′02″55	13.087	11.230	10.615	10.471	7
206011691	K7.9(1.1)	0.721(59)	3952(202)	0.737(56)	0.114(14)	22 <sup>h</sup> 41 <sup>m</sup> 12 <sup>s</sup> 885	-14°29′20″35	12.316	10.251	9.633	9.417	7
206061524	M0.7(1.2)	0.726(62)	3961(213)	0.743(59)	0.117(15)	22 <sup>h</sup> 20 <sup>m</sup> 13 <sup>s</sup> 766	-13°06′52″66	14.443	12.413	11.796	11.579	
206162305	M1.1(1.1)	0.695(58)	3896(202)	0.713(56)	0.100(13)	22h23m02s289	$-10^{\circ}29'18''89$	14.807	12.608	11.933	11.766	
206192813	M1.6(1.2)	0.622(62)	3966(225)	0.642(62)	0.086(13)	22h46 <sup>m</sup> 53 <sup>s</sup> 865	-09°52′53″83	14.875	12.598	11.927	11.732	
206209135	M2.7(0.9)	0.359(54)	3370(166)	0.361(61)	0.0149(32)	22 <sup>h</sup> 18 <sup>m</sup> 29 <sup>s</sup> 271	-09°36′44″58	14.407	11.685	11.122	10.962	7
211331236	M1.0(1.0)	0.467(66)	3781(203)	0.481(71)	0.0400(82)	08 <sup>h</sup> 55 <sup>m</sup> 25 <sup>s</sup> 364	$+10^{\circ}28'08''87$	13 905	11 447	10 801	10 589	7
211357309	M2 1(0 7)	0.506(54)	3790(179)	0.523(57)	0.0474(75)	08 <sup>h</sup> 52 <sup>m</sup> 55 <sup>s</sup> 831	$+10^{\circ}56'41''00$	13 155	10 781	10 165	9.885	7
211428897	$M_{2,7(1,2)}$	0.324(54)	3577(200)	0.321(63)	0.0155(37)	08 <sup>h</sup> 35 <sup>m</sup> 25 <sup>s</sup> 812	$+12^{\circ}04'33''04$	13 205	10 414	9 863	9 624	7
211770795	K4.3(1.3)	0.637(58)	4311(291)	0.656(58)	0.126(18)	$08^{h}48^{m}02^{s}336$	$+16^{\circ}54'06''67$	14.489	12.841	12.265	12.174	7
211799258	M3 0(0.8)	0.027(50)	3411(176)	0.257(66)	0.0089(26)	$08^{h}32^{m}59^{s}077$	$+17^{\circ}18'23''57$	15 979	13 017	12.420	12.185	, 7
211831378	M0 3(1 5)	0.608(61)	4141(250)	0.627(61)	0.098(15)	$08^{h}24^{m}33^{s}033$	$+17^{\circ}45'43''16$	16 270	13 972	13 228	13.085	
211916756	M1 2(0.9)	0.420(90)	3704(214)	0.43(10)	0.0299(93)	08 <sup>h</sup> 37 <sup>m</sup> 27 <sup>s</sup> 058	$+18^{\circ}58'36''07$	15 498	13 312	12,738	12.474	
211970234	M4 2(0.8)	0.185(57)	3000(213)	0.148(74)	0.025(11)	$09^{h}04^{m}21 \times 043$	$+10^{\circ}46'48''98$	16 122	13 851	13 186	12.987	
212006344	M0.8(0.9)	0.595(55)	3918(226)	0.615(56)	0.075(11)	08 <sup>h</sup> 25 <sup>m</sup> 54 <sup>s</sup> 315	$+20^{\circ}21'34''45$	12 466	10 104	9 4 5 7	9 275	7
212069861	M0.6(0.9)	0.692(58)	4078(223)	0.709(56)	0.119(15)	08 <sup>h</sup> 57 <sup>m</sup> 46 <sup>s</sup> 605	$+21^{\circ}27'12''72$	14 102	11 907	11 250	11.055	7
212154564	M2 7(1 2)	0.32(15)	3502(162)	0.32(18)	0.0140(94)	08 <sup>h</sup> 54 <sup>m</sup> 33 <sup>s</sup> 884	$+23^{\circ}07'58''40$	15 105	12.838	12 227	11.035	7
212315941	K7 9(0 9)	0.32(13) 0.48(13)	4056(219)	0.32(10) 0.49(14)	0.057(22)	13 <sup>h</sup> 32 <sup>m</sup> 20 <sup>s</sup> 944	$-17^{\circ}03'40''29$	14 406	12.830	12.227	12 175	
212354731	M1 8(1 1)	0.356(55)	3369(166)	0.356(62)	0.0147(33)	13 <sup>h</sup> 33 <sup>m</sup> 22 <sup>s</sup> 379	-16°00/23"85	15 805	13 412	12.200	12.507	7
212565386	M1 0(0.8)	0.550(55)	3989(228)	0.590(58)	0.074(11)	$13^{h}30^{m}26^{s}554$	$-11^{\circ}20'29''42$	14 727	12.368	11 746	11 513	, 7
212679798	M0.6(1.1)	0.545(53)	3716(171)	0.563(55)	0.0508(74)	13 <sup>h</sup> 29 <sup>m</sup> 56 <sup>s</sup> 550	-08°44′58″70	14 846	13.056	12.469	12.338	, 7
212756297	K4 6(0.8)	0.717(60)	4242(257)	0.735(58)	0.150(20)	$13^{h}50^{m}37^{s}408$	-06°48′14″42	13 009	11 350	10 794	10.619	
212773309	M0.6(0.8)	0.506(56)	3886(199)	0.523(59)	0.0524(86)	$13^{h}49^{m}32^{s}380$	-06°19′21″87	11.391	9.802	9.272	9.114	7
		0.000(00)	2000(177)	0.020(0))	0.002 .(00)	10 10 021000	00 17 211 07		2.001	··-·-	···· ·	

 $\neg$ 

<sup>a</sup> Stars that are not Gliese stars (GJ) are the EPIC ID of K2 stars.

<sup>b</sup> Notes indicate stars with interferometrically determined radii and temperatures from (1) von Braun et al. (2014), (2) Boyajian et al. (2012b), (3) von Braun et al. (2012), (4) Demory et al. (2009), (5) von Braun et al. (2011), and (6) Boyajian et al.

(2008). (7) indicates those stars with parameters reported in the companion paper by Dressing et al. (2017). Finally, (8) indicates the averages of spectra obtained between two separate nights.

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Figure 4. We show stellar radius and effective temperature for all our K2 target stars (black points with error bars) derived using Equations (2) and (3). The red squares and dashed line show the average values for each SpT as calculated by Boyajian et al. (2012b).



Figure 5. Comparison of our stellar radii to those of Dressing et al. (2017). The dotted line shows a 1:1 agreement, while any deviation from the dotted line presents the small discrepancies. Overall, there is a general agreement between our works in deriving our stellar radii.

have compositions dominated by rock or iron, while larger planets are more likely to be volatile-rich. However, there may still be rocky planets larger than this limit. For example, Buchhave et al. (2016) found that *Kepler*-20b, a 1.9  $R_{\oplus}$  planet, has a density consistent with a rocky composition even though it is beyond the rocky-to-gaseous transition.

We compare our calculations of the insolation flux from our K2 sample to those from Crossfield et al. (2016) in Figure 10. The discrepancies between our values and those in Crossfield et al. (2016) highlight the importance of using spectroscopically derived stellar parameters in order to compute planet parameters.

## 5. Conclusion and Future Prospects

In this paper, we derive stellar and planetary parameters for K2 K and M dwarf systems. We adopt similar calibration techniques from Neves et al. (2014), Maldonado et al. (2015), and Terrien et al. (2015) by measuring EWs in the near-infrared part of the spectrum. Interferometric calibration samples are



**Figure 6.** Comparison of our effective temperatures to those of Dressing et al. (2017). The dotted line shows a 1:1 agreement, while any deviation from the dotted line presents the small discrepancies. Overall, there is a general agreement between our works in deriving our effective temperatures. We address small caveats in Section 4.3 for the outlier, EPIC 211770795.



**Figure 7.** Comparison of our stellar radii to those of Huber et al. (2016). The dotted line shows a 1:1 agreement, while any deviation from the dotted line presents the small discrepancies. As discussed in Section 4.3, we find that the majority of the objects in the sample are larger in our work and find a 0.15  $R_{\odot}$  median increase.



**Figure 8.** Comparison of our effective temperatures to those of Huber et al. (2016). The dotted line shows a 1:1 agreement, while any deviation from the dotted line presents the small discrepancies. See Section 4.3 for a discussion.

 Table 5

 K2 Planet and Candidate Parameters

Name <sup>a</sup>	EPIC	Р	<i>R</i> *	T <sub>eff</sub>	<i>M</i> *	a	Sinc	R <sub>P</sub>	T <sub>eq</sub>
		(days)	$(R_{\odot})$	(K)	$(M_{\odot})$	(au)	$(S_{\oplus})$	$(R_{\oplus})$	(K)
K2-43b	201205469.01	3.471140	0.559(57)	3923(198)	0.577(58)	0.0374(13)	47.7(8.2)	4.01(45)	720
K2-4b	201208431.01	10.004438	0.658(56)	3900(195)	0.678(56)	0.0800(22)	14.0(2.0)	2.52(37)	530
K2-3b	201367065.01	10.054428	0.565(61)	3976(205)	0.584(62)	0.0762(27)	12.4(2.2)	2.15(26)	510
K2-3c	201367065.02	24.643479	0.565(61)	3976(205)	0.584(62)	0.1384(50)	3.72(68)	1.76(22)	380
K2-3d	201367065.03	44.560906	0.565(61)	3976(205)	0.584(62)	0.2055(75)	1.70(30)	1.44(18)	310
K2-9b	201465501.01	18.447385	0.366(53)	3460(164)	0.370(61)	0.0980(56)	1.78(45)	4.9(1.1)	320
	201617985.01	7.281384	0.608(55)	3868(193)	0.627(56)	0.0630(19)	19.0(2.9)	27(23)	570
K2-49b	201690311.01	2.770645	0.697(94)	3948(203)	0.713(92)	0.0346(15)	89(18)	2.90(44)	840
	201717274.01	3.527432	0.368(80)	3528(165)	0.371(93)	0.0326(29)	18.1(6.8)	1.55(39)	560
K2-18b	201912552.01	32.941798	0.411(53)	3527(162)	0.419(58)	0.1502(70)	1.04(23)	2.31(31)	280
	204489514.01	10.223626	0.230(56)	3096(198)	0.206(71)	0.0544(66)	1.5(1.1)	14.8(6.6)	300
K2-54b	205916793.01	9.784339	0.707(60)	4103(230)	0.725(57)	0.0804(21)	19.8(2.8)	2.10(27)	580
K2-55b	205924614.01	2.849258	0.769(63)	4240(259)	0.784(59)	0.03620(90)	131(18)	4.63(40)	920
K2-21b	206011691.01	9.323890	0.721(59)	3952(202)	0.739(57)	0.0786(20)	18.4(2.5)	1.92(18)	560
K2-21c	206011691.02	15.501158	0.721(59)	3952(202)	0.739(57)	0.1101(28)	9.4(1.3)	2.37(24)	480
	206061524.01	5.879750	0.726(62)	3961(213)	0.742(58)	0.0576(15)	35.0(5.2)	6.92(61)	660
K2-69b	206162305.01	7.065991	0.695(58)	3896(202)	0.714(56)	0.0644(17)	24.1(3.3)	3.25(37)	600
K2-71b	206192813.01	6.985406	0.622(62)	3966(225)	0.640(62)	0.0615(20)	22.9(3.7)	3.11(42)	600
K2-72b	206209135.01	5.577387	0.359(54)	3370(166)	0.362(61)	0.0438(26)	7.8(2.0)	1.15(20)	460
K2-72c	206209135.02	15.187114	0.359(54)	3370(166)	0.362(61)	0.0855(50)	2.05(53)	1.30(22)	330
K2-72d	206209135.03	7.759932	0.359(54)	3370(166)	0.362(61)	0.0548(32)	5.0(1.3)	1.10(21)	410
K2-72e	206209135.04	24.166851	0.359(54)	3370(166)	0.362(61)	0.1163(67)	1.11(28)	1.25(24)	280
	211331236.01	1.291651	0.467(66)	3781(203)	0.481(71)	0.01819(92)	119(28)	1.46(57)	900
	211428897.01	1.610918	0.324(54)	3577(200)	0.320(64)	0.0183(13)	46(13)	0.86(16)	710
	211770795.01	7.729341	0.637(58)	4311(291)	0.656(57)	0.0665(19)	28.2(4.5)	2.3(1.3)	630
	211799258.01	19.535120	0.271(54)	3411(176)	0.257(66)	0.0900(82)	1.12(41)	15.7(7.9)	280
K2-95b	211916756.01	10.133866	0.420(90)	3704(214)	0.43(10)	0.0691(60)	6.2(2.4)	13(14)	430
	211970234.01	1.483459	0.185(57)	3000(213)	0.150(75)	0.0136(23)	13(35)	10.0(6.8)	520
	212006344.01	2.219215	0.595(55)	3918(226)	0.615(56)	0.02835(87)	93(14)	1.28(15)	850
	212069861.01	30.953052	0.692(58)	4078(223)	0.709(56)	0.1721(44)	4.01(57)	3.12(36)	390
	212154564.01	6.413647	0.32(15)	3502(162)	0.32(18)	0.0464(91)	7.1(0.3)	2.5(1.2)	450
	212315941.01	12.935695	0.48(13)	4056(219)	0.50(14)	0.0851(87)	8.0(4.1)	5.9(3.7)	460
	212354731.01	20.397357	0.356(55)	3369(166)	0.358(63)	0.1037(60)	1.34(36)	25.5(9.7)	290
	212679798.01	1.834810	0.545(53)	3716(171)	0.562(55)	0.02417(80)	86(13)	28(18)	830
	212756297.01	1.337116	0.717(60)	4242(257)	0.733(58)	0.02142(56)	326(45)	13.7(1.2)	1160

Note.

<sup>a</sup> K2 names indicate validated planets, while those without a K2 name indication remain planet candidates.

used from Demory et al. (2009), von Braun et al. (2011, 2012), Boyajian et al. (2012b), and von Braun et al. (2014) in order to provide a more precise baseline to calculate the stellar radii and effective temperatures of the stars in our sample. Various functions (whether they are linear, quadratic, or a ratio of EWs) are tested, and we use the functions with the best BIC value and the lowest residuals to calculate stellar parameters.

Our spectroscopically derived stellar radii improve on previously reported values that relied on stellar models poorly calibrated to these low-mass stars. We find a median increase of  $0.15 R_{\odot}$  when comparing our measurements to those of Huber et al. (2016), consistent with the median increase in size found by Newton et al. (2015) when revising the photometrically based stellar radius estimates determined by Dressing & Charbonneau (2013) for cool dwarfs observed during the prime *Kepler* mission. Finally, we calculate the *K2* planet or planet candidate radius and equilibrium temperature.

Since our team also obtained optical spectra, using the EFOSC2 spectrograph (Buzzoni et al. 1984) on the NTT, in a future work we will apply the same techniques in order to cross-check our stellar properties. Furthermore, this work does

not calculate stellar metallicities; however, we plan to do so in later works.

Our work paves the way for future exoplanet surveys. Other spectroscopic and photometric surveys focusing on M dwarfs are currently underway or are being planned for the near future. SPECULOOS, a 1 m near-infrared telescope, will observe approximately 500 of the nearest M and brown dwarfs in the southern hemisphere (Gillon et al. 2013). CARMENES will provide high-resolution (R = 82,000) spectra between 0.5 and 1.7  $\mu$ m for late-type M dwarfs and search for Earth-like planets in the habitable zone (Quirrenbach et al. 2012). The Habitable Zone Planet Finder will also provide spectra for M dwarfs and will attempt to find planets through the Doppler effect (Mahadevan et al. 2010). Yet another RV survey, SPIRou, aims to find exoplanets around low-mass stars using high-resolution spectra between 0.98 and 2.35  $\mu$ m (Santerne et al. 2013).

Future transit surveys will detect many new Earth-like planets around M dwarfs, just like previous and ongoing photometric surveys such as *Kepler* and *K2*. Although the current *Gaia* mission (Lindegren 2010) focuses more on astrometry (for which stellar mass is a key input), its two



**Figure 9.** Planet radii, incident irradiation, and equilibrium temperatures of all K2 planets and candidates observed in our program. Venus and Earth are indicated by single letters. Red plus signs indicate validated planets, and gray squares indicate planet candidates, as reported by Crossfield et al. (2016). The shaded region represents the approximate location of the cloud-free habitable zone for an early-type M dwarf (Kopparapu et al. 2013). That zone was defined for planets with masses 0.3–10 times that of Earth. The larger of those masses corresponds to the upper, lightly shaded area (Wolfgang et al. 2016).



**Figure 10.** Planet radii and incident irradiation for the *K2* planets and planet candidates that appear in both our work and Crossfield et al. (2016). Black circles indicate the *K2* objects reported by Crossfield et al. (2016), while the red squares indicate the *K2* objects in this work. The gray lines connect our updated parameter values to the original estimates published by Crossfield et al. (2016). See Section 4.4 for details.

photometers can provide light curves for exoplanet detection. The *Transiting Exoplanet Survey Satellite (TESS*; Ricker et al. 2009) and *PLAnetary Transits and Oscillations of stars (PLATO*; Rauer et al. 2014) will also find planets, some of which will be high-priority targets for the *James Webb Space Telescope (JWST*; Gardner et al. 2006). The recent announcement of a roughly Earth-mass planet candidate orbiting Proxima Centauri (Anglada-Escudé et al. 2016) adds yet more urgency to the need to search for more planets and characterize their low-mass host stars. The combination of all of these surveys will yield many new M dwarf systems in need of stellar

and planetary parameters and of a large, precise calibration sample.

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