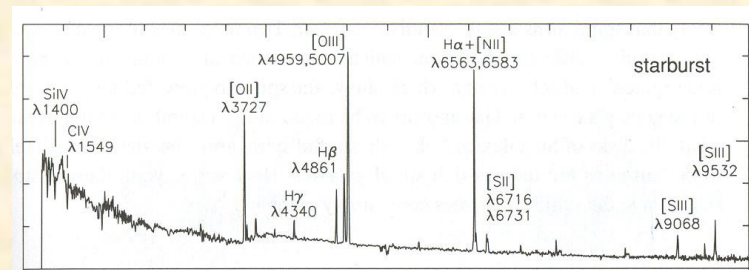
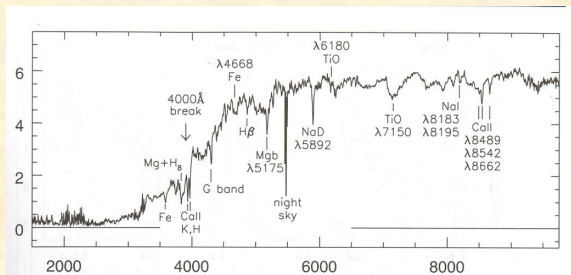
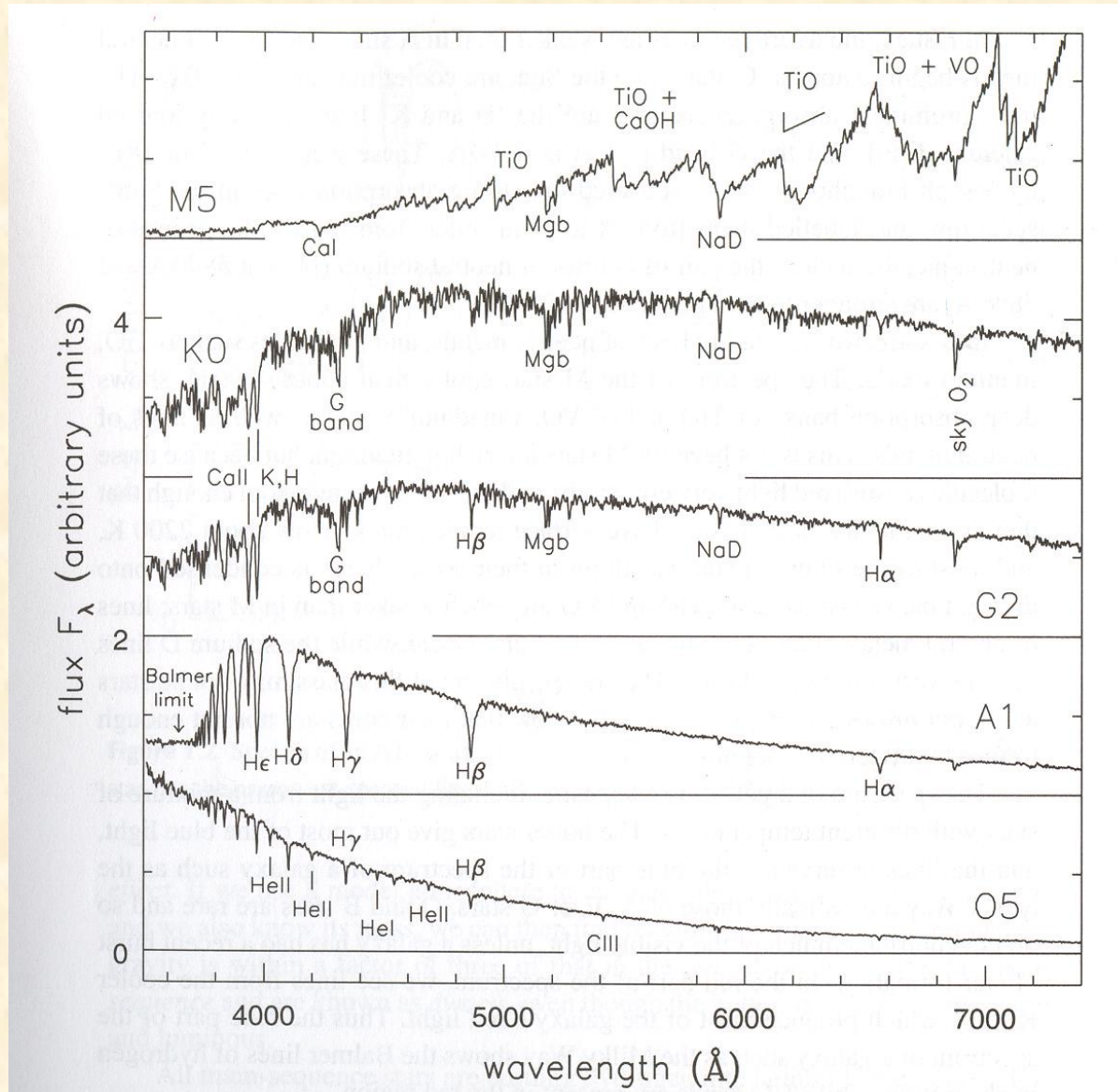


Kinematics of Galaxies

- Spectral Features of Galaxies
- Basics of Spectroscopy
- Elliptical Kinematics
- Faber-Jackson and the Fundamental Plane
- Disk Kinematics (Stellar and H I)
- 2D Velocity Fields
- Rotation Curves and Masses
- Tully-Fisher
- Detection of Supermassive Black Holes

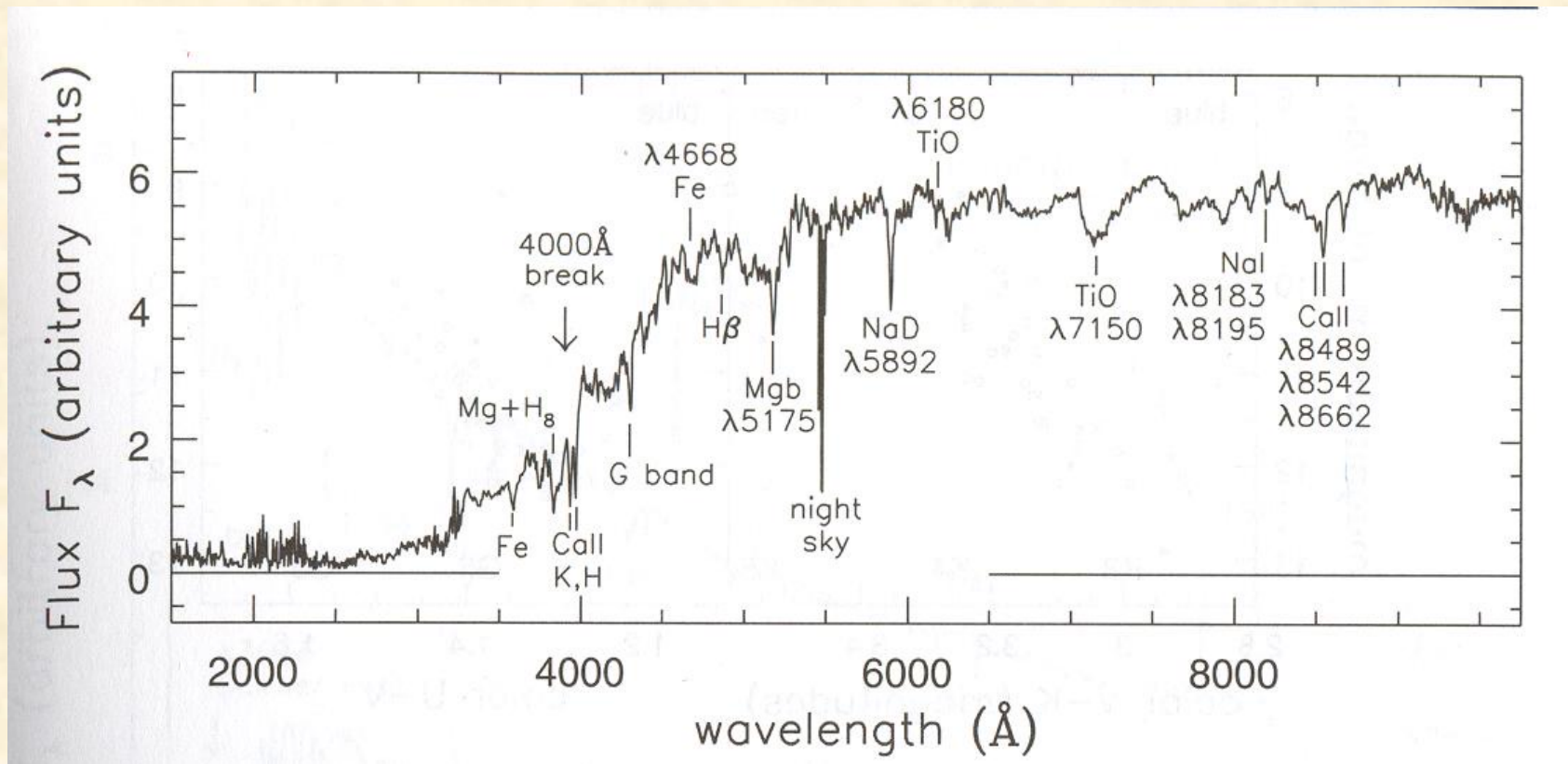


Stellar Spectra



(Sparke and Gallagher, p. 5)

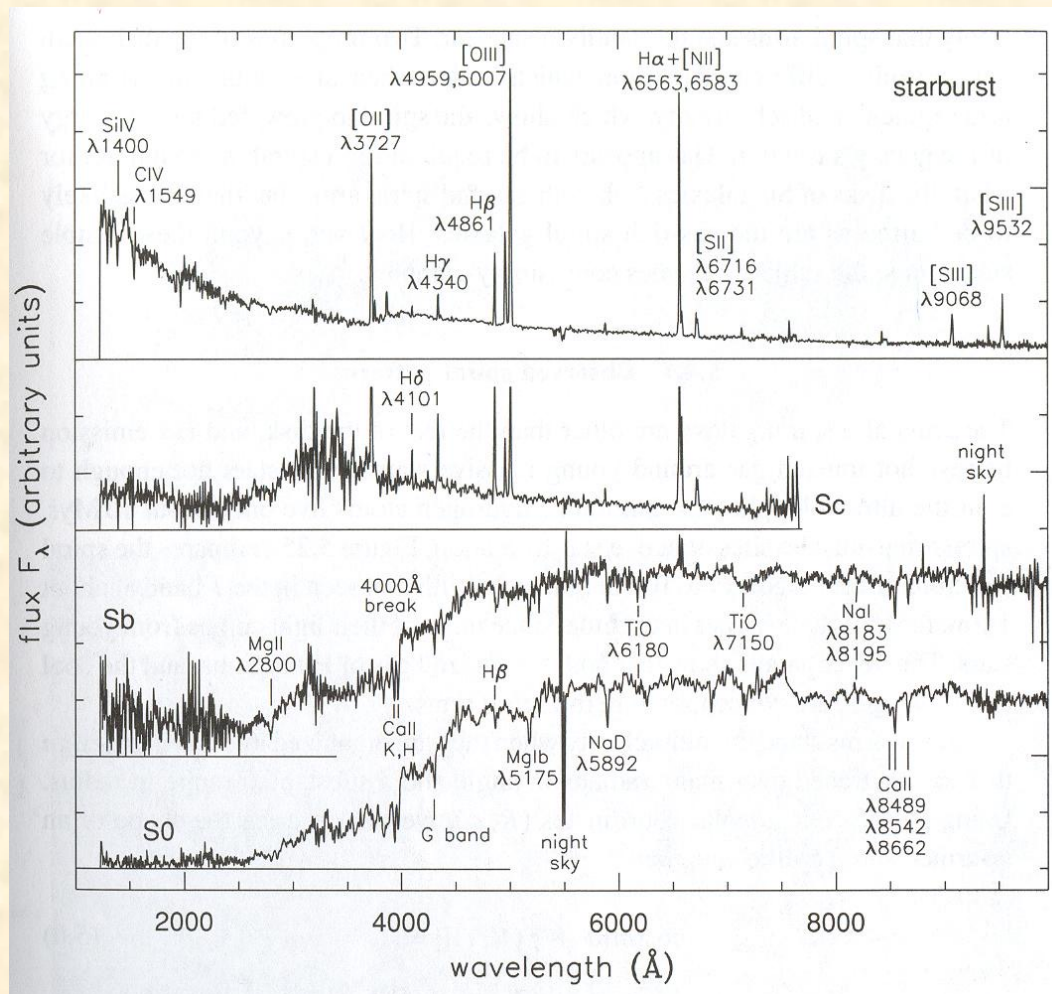
Galaxy Spectra - Ellipticals



(Sparke and Gallagher, p. 267)

- Most features from giant G and K stars (e.g., G band is from CH)
- In the optical, most absorption is stellar. Ca II H, K and Na I D can come from ISM as well (but not much in Ellipticals)
- Lines are broadened from stellar motions
- Ca II triplet lines at $\sim 8500 \text{ \AA}$ are good for kinematics (well separated, uncontaminated)

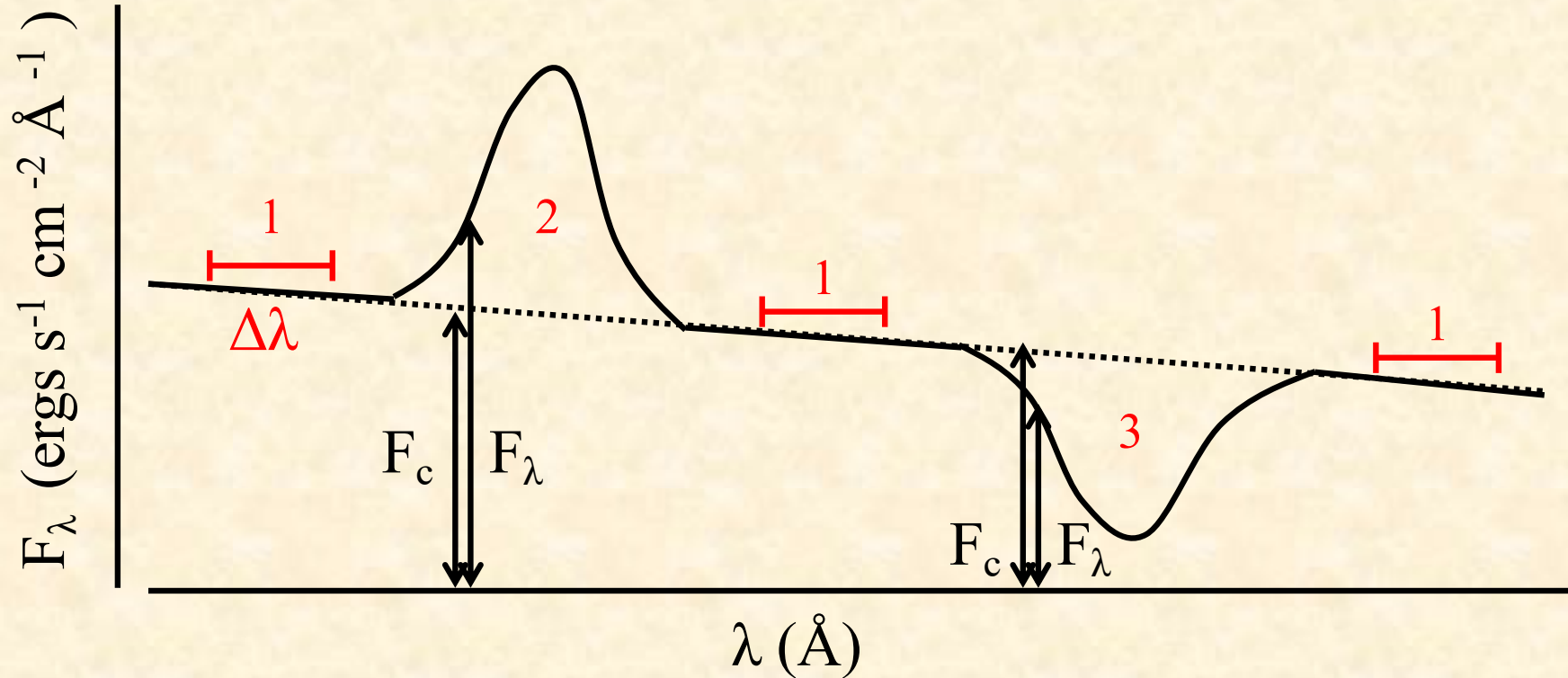
Disk Galaxies



(Sparke and
Gallagher, p. 224)

- S0 similar to E' s \rightarrow old stellar populations
- Sa/Sb have stronger Balmer lines (A, F stars) and bluer continua
- Sc have emission lines from H II regions (young hot stars)
- Starburst galaxies have very strong emission lines and blue continua

Basics of Spectroscopy

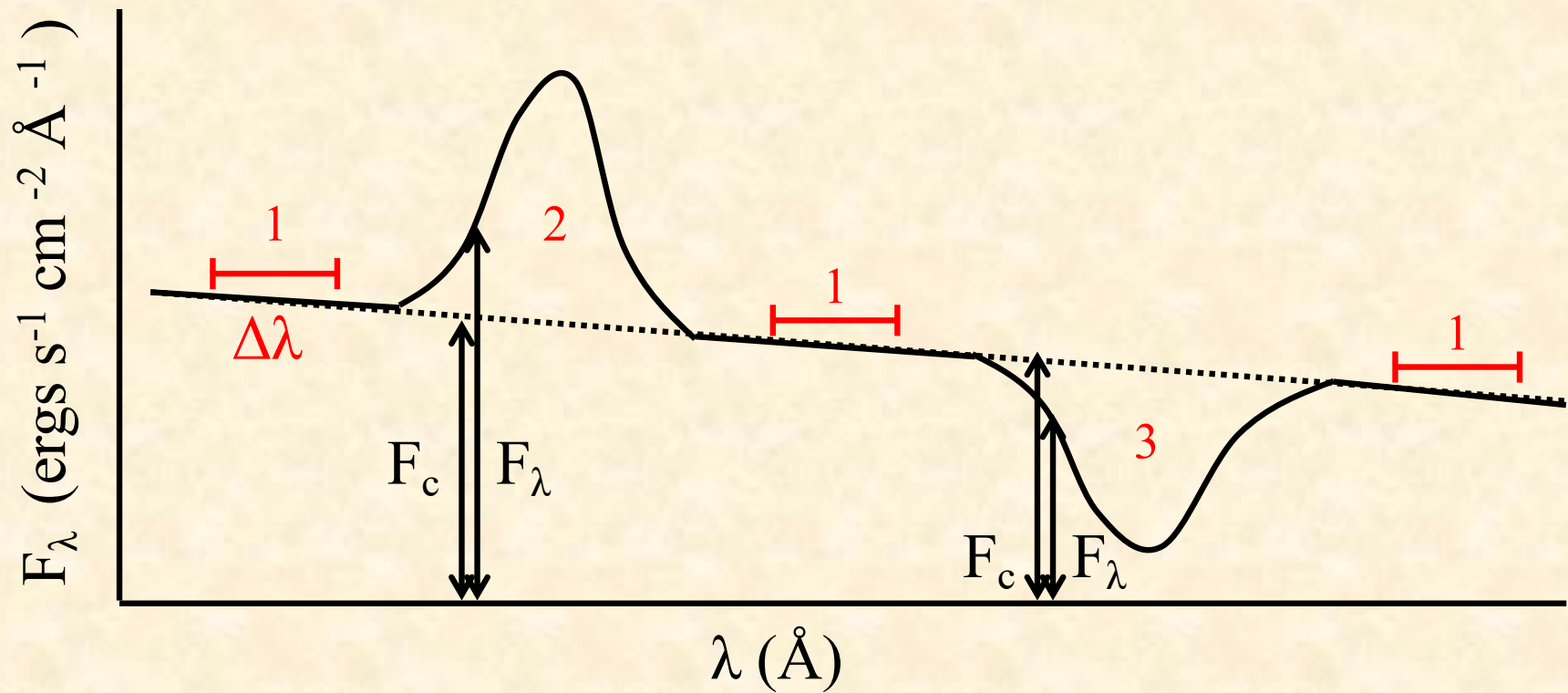


Units:

1. Continuum Flux: $F_c = \frac{\int F_\lambda d\lambda}{\Delta\lambda} = \langle F_\lambda \rangle$ (ergs s⁻¹ cm⁻² Å⁻¹)

2. Emission Line Flux: $F = \int (F_\lambda - F_c) d\lambda$ (ergs s⁻¹ cm⁻²)

3a. Absorption Equivalent Width: $W_\lambda = \int (1 - F_\lambda/F_c) d\lambda$ (Å)



3b. Absorption-Line Centroid:

$$\lambda_c = \frac{\int \lambda (F_c - F_\lambda) d\lambda}{\int (F_c - F_\lambda) d\lambda} \quad (\text{\AA})$$

3c. Radial Velocity Centroid:
(nonrelativistic)

$$v_r = \frac{\lambda_c - \lambda_{\text{lab}}}{\lambda_{\text{lab}}} c \quad (\text{km s}^{-1})$$

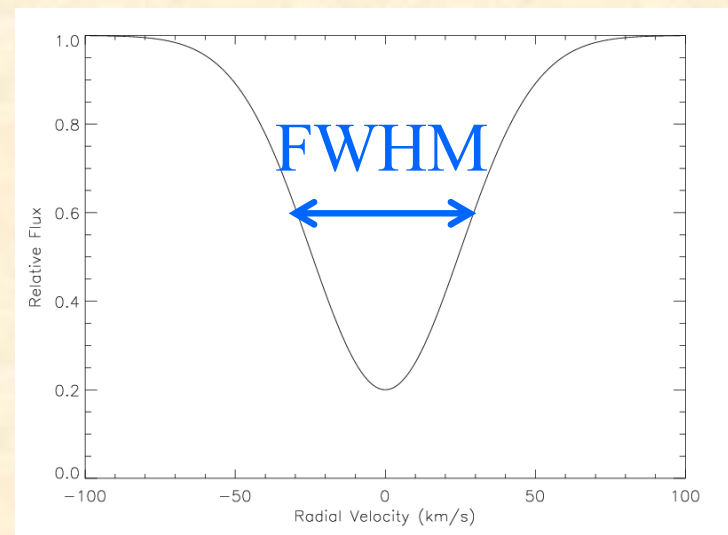
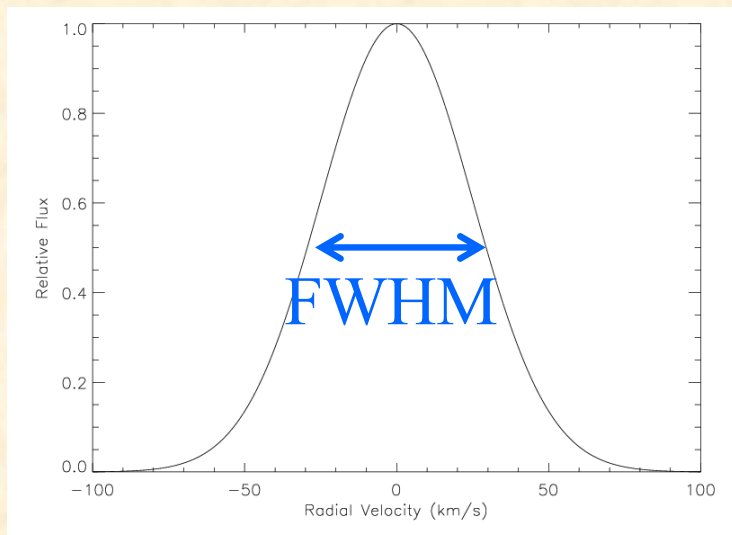
- For Galactic kinematics, v_r and σ are used
- A **Gaussian** profile is often assumed for the LOSVD (line of sight velocity distribution):

$$P(v_r) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{1}{2}(v_r/\sigma)^2}$$

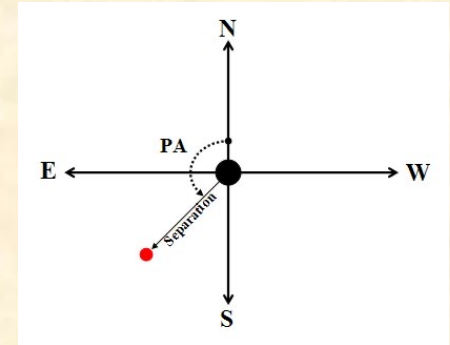
where v_r = centroid = peak, σ = velocity dispersion

- Note the full-width at half-maximum for a Gaussian is:

$$\text{FWHM} = 2.355 \sigma$$



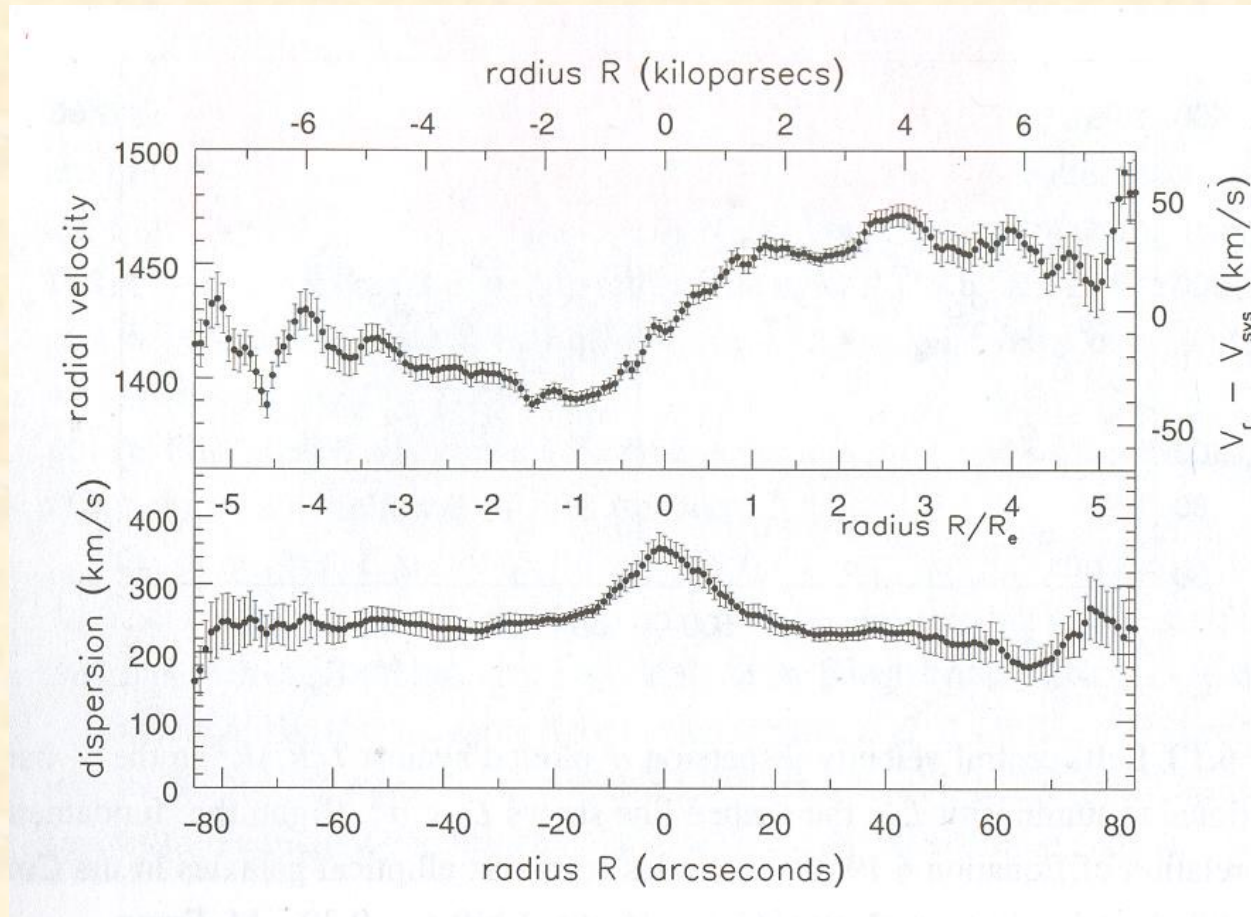
Spatially-Resolved Spectra



- Long-slit spectroscopy: spectra at each position along slit
- Resolving power needed: $R = \lambda/\Delta\lambda \approx 5000$
(where $\Delta\lambda$ is the FWHM of the line-spread function (LSF))
- Measure v_r and σ at each position.
- Subtract **systemic** velocity (due to Hubble flow, etc.) from v_r
- Net v_r at each position is a measure of rotation: $v_r = v \sin(\text{incl})$
- σ gives component of random motion in the line of sight

Ellipticals: Kinematics

Ex) cD galaxy NGC 1399



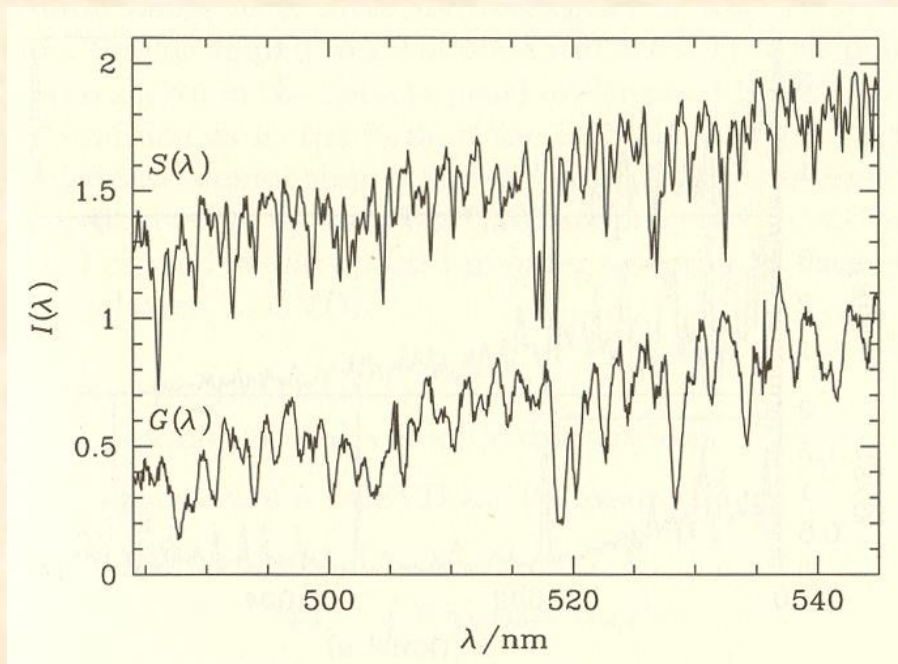
(Sparke and Gallagher, p. 257)

- For most E' s: v_r (max) $\ll \sigma$ (central velocity dispersion)

Determination of v_r and σ

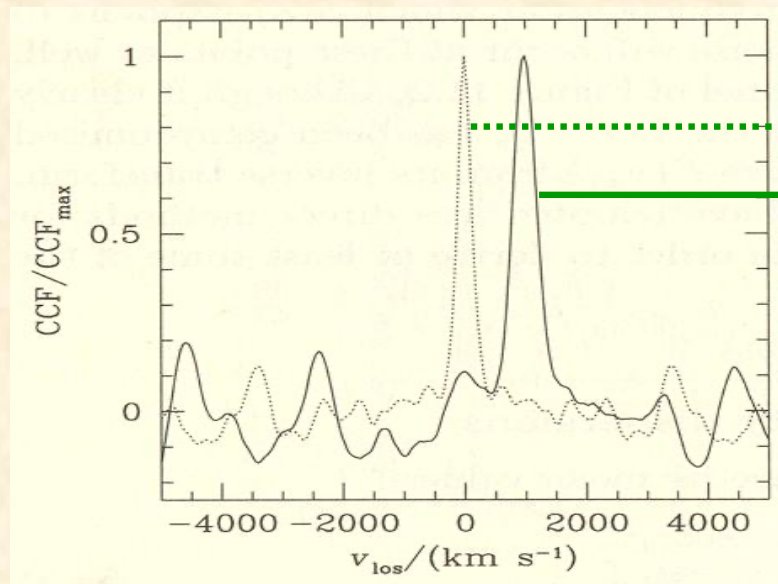
- One method: use cross-correlation function (CCF)
- Cross-correlate the galaxy spectrum with that of a star (like a K giant) or a synthetic galaxy
 - At each λ , you have $F_\lambda(\text{star})$ and $F_\lambda(\text{galaxy})$
 - Do a linear fit of $F_\lambda(\text{galaxy})$ vs. $F_\lambda(\text{star})$ to get “r” (linear-correlation coefficient) (Bevington, p. 121)
 - Shift one spectrum in λ , and calculate r again
($r = 1 \rightarrow$ perfect correlation; $r = 0 \rightarrow$ no correlation)
 - The CCF is just r as a function of shift
- The CCF peak give the velocity centroid v_r ; the CCF width gives σ
- The auto-correlation function (ACF) is a function cross-correlated with itself.

CCF Example



→ K0 giant

→ NGC 2549



(Binney & Merrifield p.695, 698)

Results for Ellipticals: Kinematic Correlations

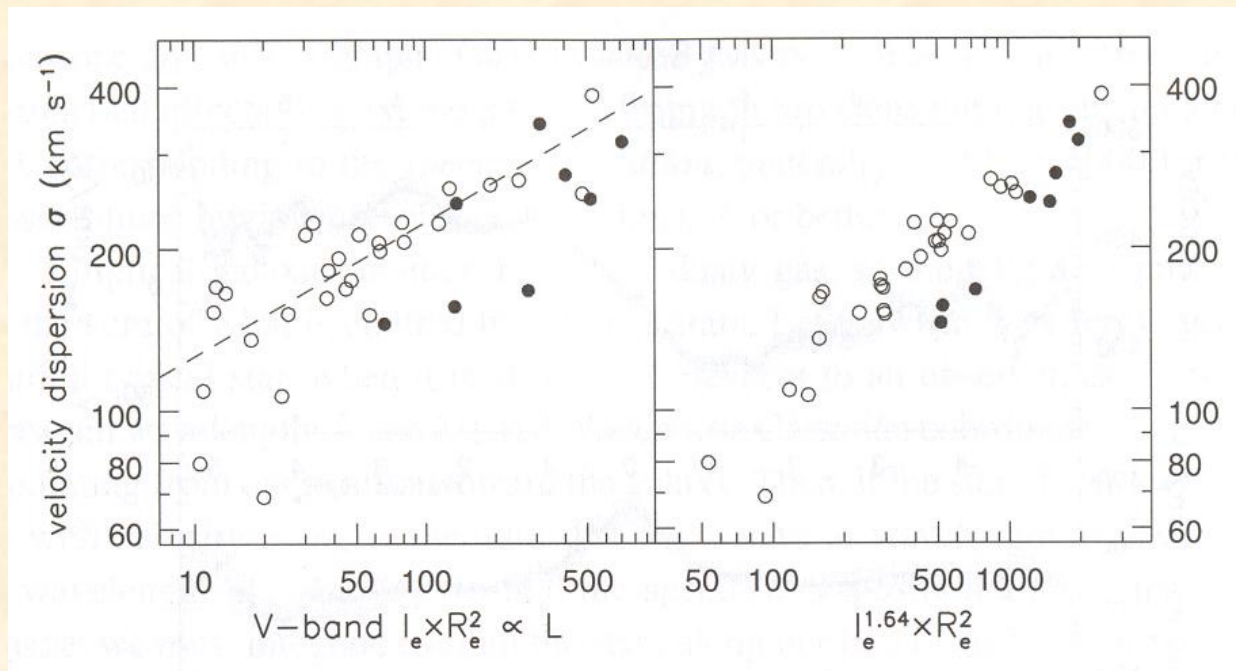
- Faber-Jackson relation: $L \sim \sigma^4$ (σ : central velocity disp.)

$$\frac{L_V}{2 \times 10^{10} L_{\odot}} = \left(\frac{\sigma}{200 \text{ km s}^{-1}} \right)^4$$

- Note $L \sim I_e R_e^2 \rightarrow$ Is there a tighter relationship for σ , I_e , R_e ?

Faber-Jackson

Fundamental Plane



(Sparke and Gallagher, p. 258)

\rightarrow projection:

$$I_e^{1.64} R_e^2 \propto \sigma^{2.48}$$

Note: These relations do *not* apply to diffuse E's and dwarf spheroidals

Rotation of Elliptical Galaxies

- Is the oblateness of ellipticals due to rotation?
 - no, E' s tend to rotate more slowly than they should
- How fast should they rotate?
- Virial Theorem – if the galaxy is dynamically relaxed, velocity dispersions are equal in all directions and:

$$2\langle KE_i \rangle + \langle PE_i \rangle = 0 \quad \text{for } i = x, y, z \text{ (axes of symmetry)}$$

where $\langle PE_i \rangle$ is the average gravitational potential.

For an oblate galaxy rotating around the z axis:

$$\frac{\langle PE_z \rangle}{\langle PE_x \rangle} = \frac{\langle KE_z \rangle}{\langle KE_x \rangle} = \frac{\sigma_z^2}{\frac{1}{2}v^2 + \sigma_x^2}$$

$$\frac{\langle PE_z \rangle}{\langle PE_x \rangle} \approx (B / A)^{0.9} = (1 - e)^{0.9} \quad (\text{Sparke \& Gallagher, 260})$$

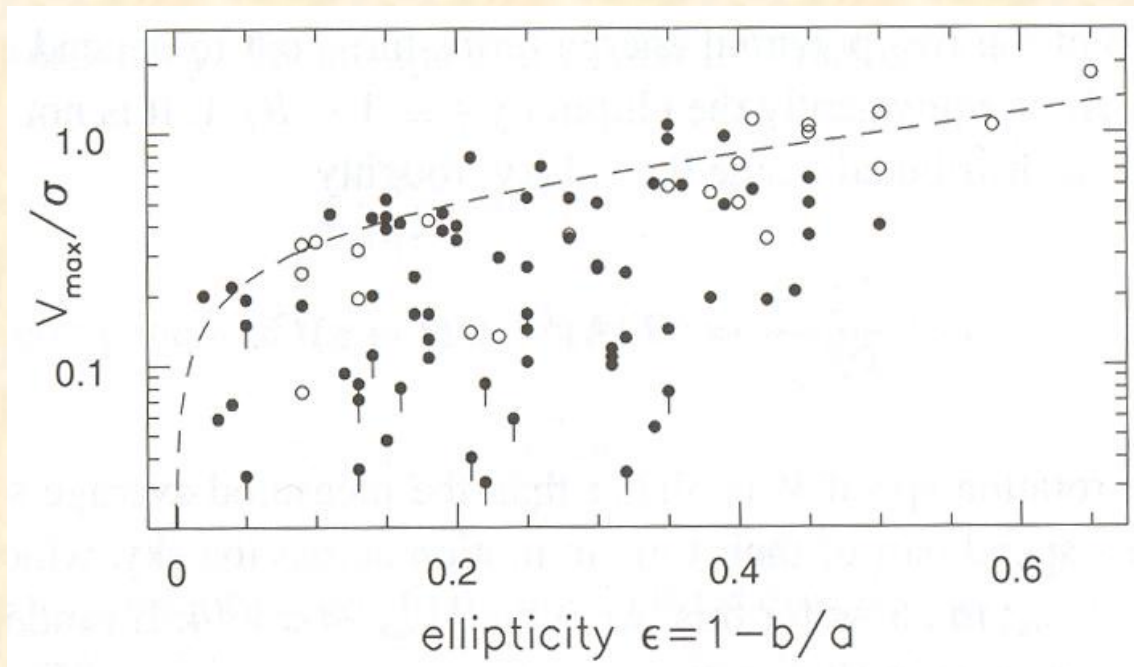
where A, B, and e are the actual axes and ellipticity

If the virial theorem applies, $\sigma = \sigma_x = \sigma_y = \sigma_z$

Note that the maximum radial velocity at σ ($= \sigma_0$) is:

$$v(\text{max}) \approx \frac{\pi}{4} v \quad (\text{Sparke \& Gallagher, p. 260})$$

$$\text{Thus } \frac{v(\text{max})}{\sigma} = \frac{\pi}{4} \sqrt{2 \left[(1-e)^{-0.9} - 1 \right]}$$

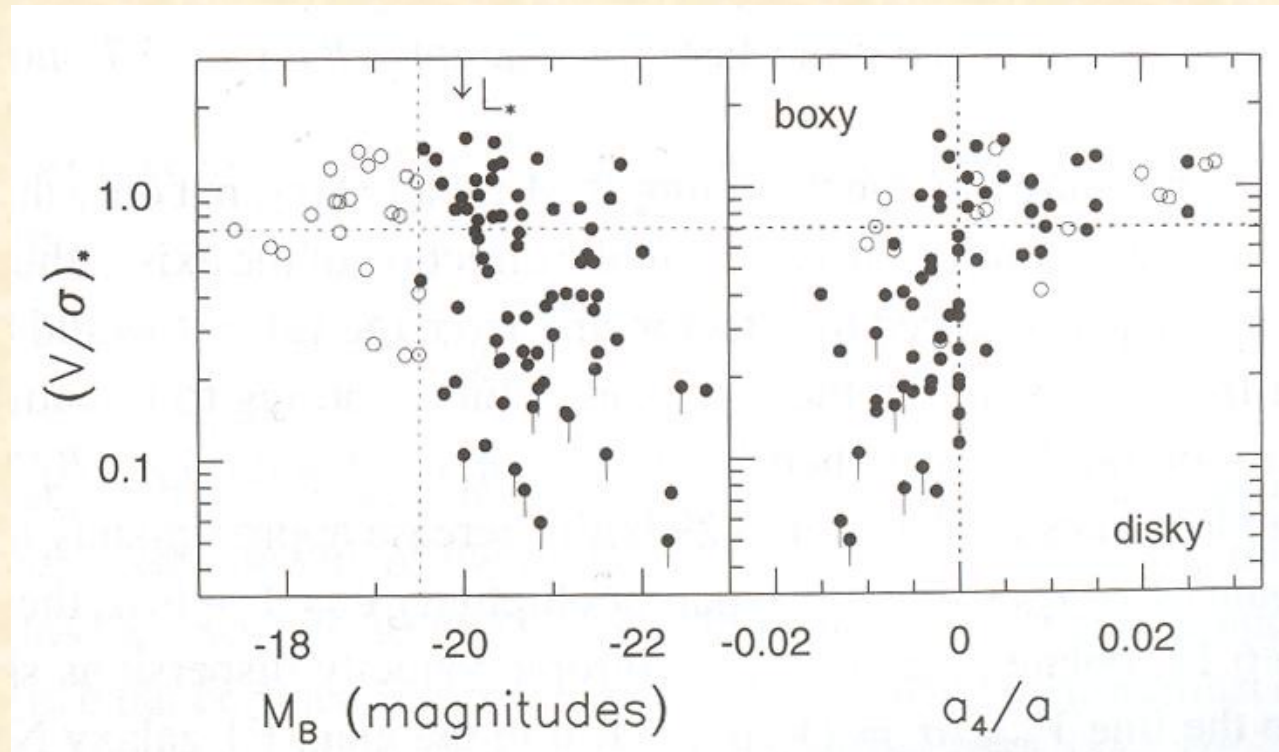


(Sparke & Gallagher,
p. 262)

→ observed $v(\text{max})$ much lower than expected from relaxed systems

- So most ellipticals are not supported by rotation, but by anisotropic velocity dispersions: $\sigma_x \neq \sigma_y \neq \sigma_z$.

Observations: Let $(v/\sigma)_* = \frac{(v_{\max}/\sigma_0)_{\text{obs}}}{(v_{\max}/\sigma_0)_{\text{eqn}}}$

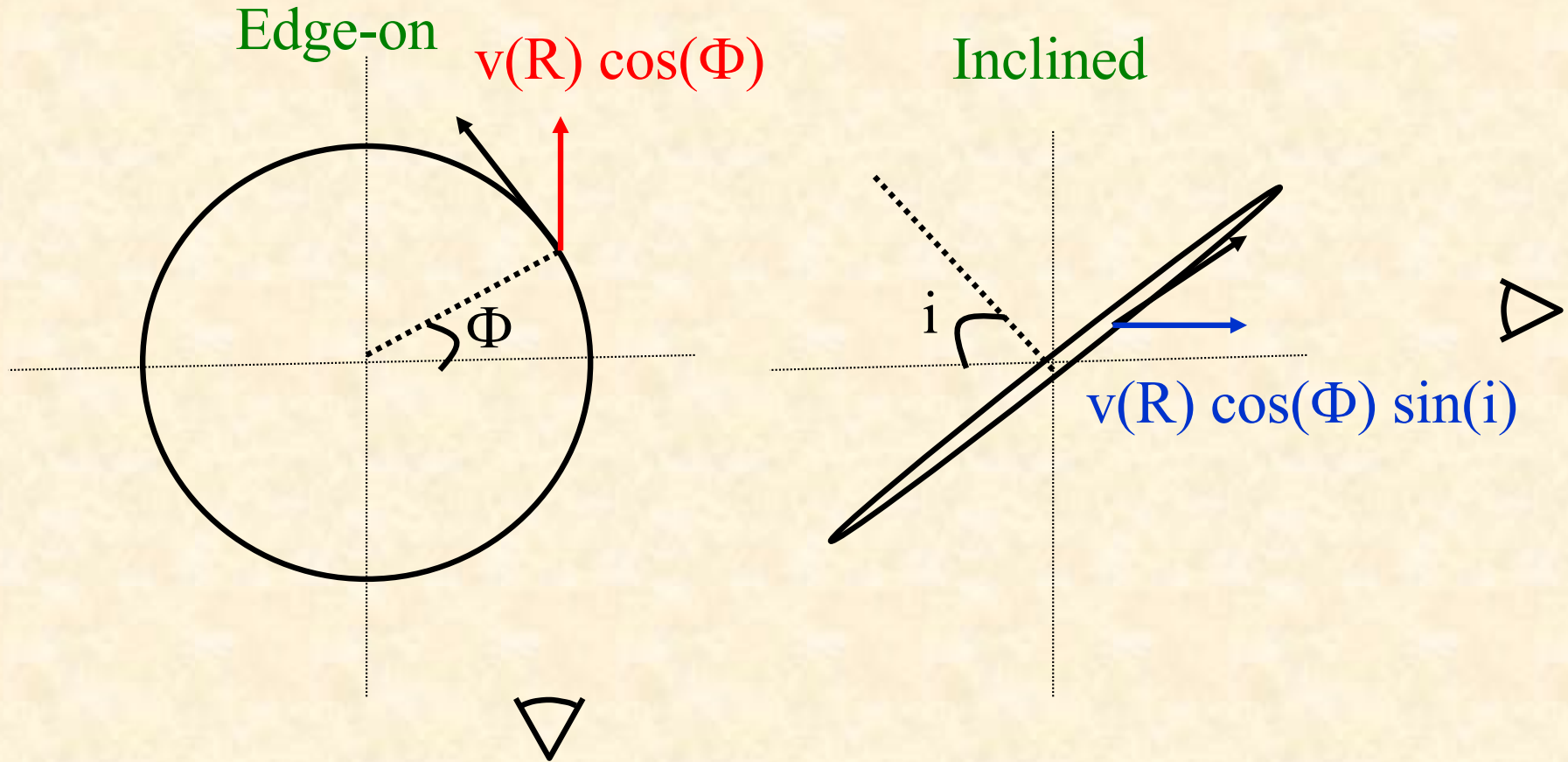


(Sparke & Gallagher, p. 261)

- Luminous and boxy ellipticals rotate much slower than expected.
- Disky E' s may be composite: rotating disk embedded in normal elliptical

Kinematics of Rotating Disks (Spirals)

- Spiral galaxies are dominated by rotation ($v_r \geq 10\sigma$)
- Can determine true velocity $v(R)$, since we know inclination



- Observed radial velocity: $v_r = v_{\text{sys}} + v(R) \cos(\Phi) \sin(i)$

Axisymmetric Disk:

$$v^2(R) = \frac{GM(r < R)}{R}$$

1) If $M \sim R$: $v = \text{const.}$

$$V_r - V_{\text{sys}} \sim \cos(\Phi) \sin(i)$$

2) If $M \sim R^3$: $v \sim R$

$$V_r - V_{\text{sys}} \sim R \cos(\Phi) \sin(i)$$

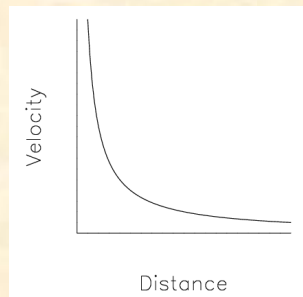
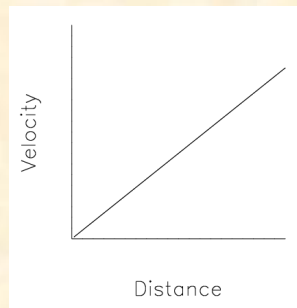
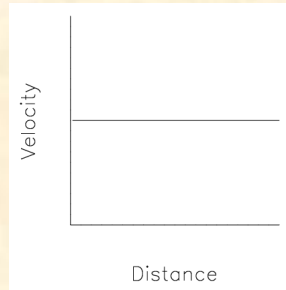
(constant density)

3) If $M = \text{constant}$: $v \sim R^{-1/2}$

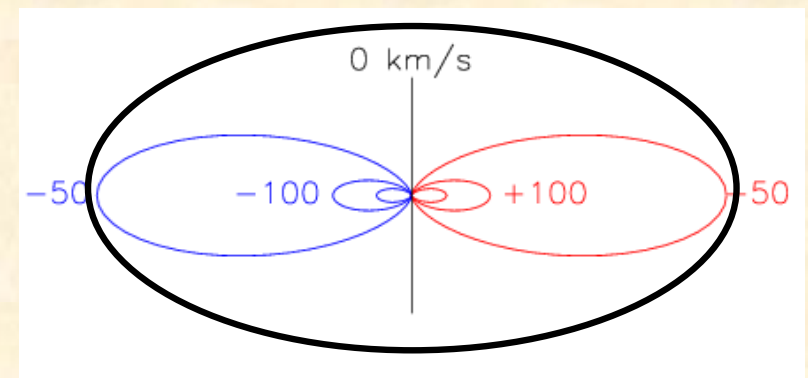
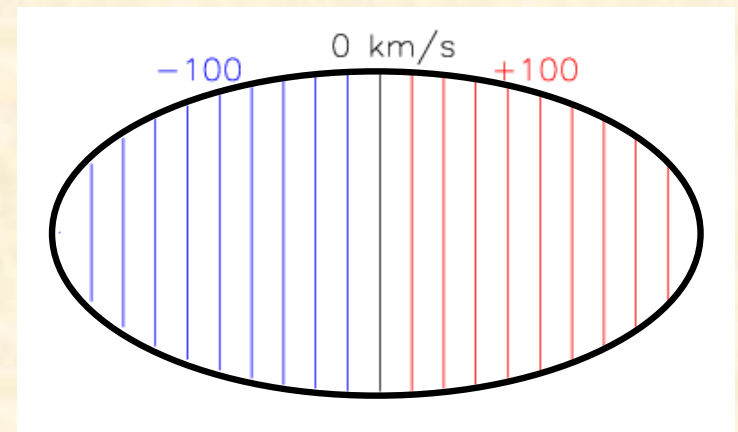
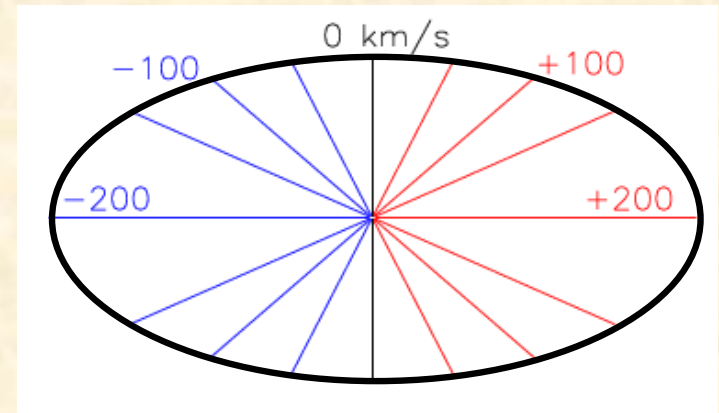
$$V_r - V_{\text{sys}} \sim R^{-1/2} \cos(\Phi) \sin(i)$$

(point-source mass)

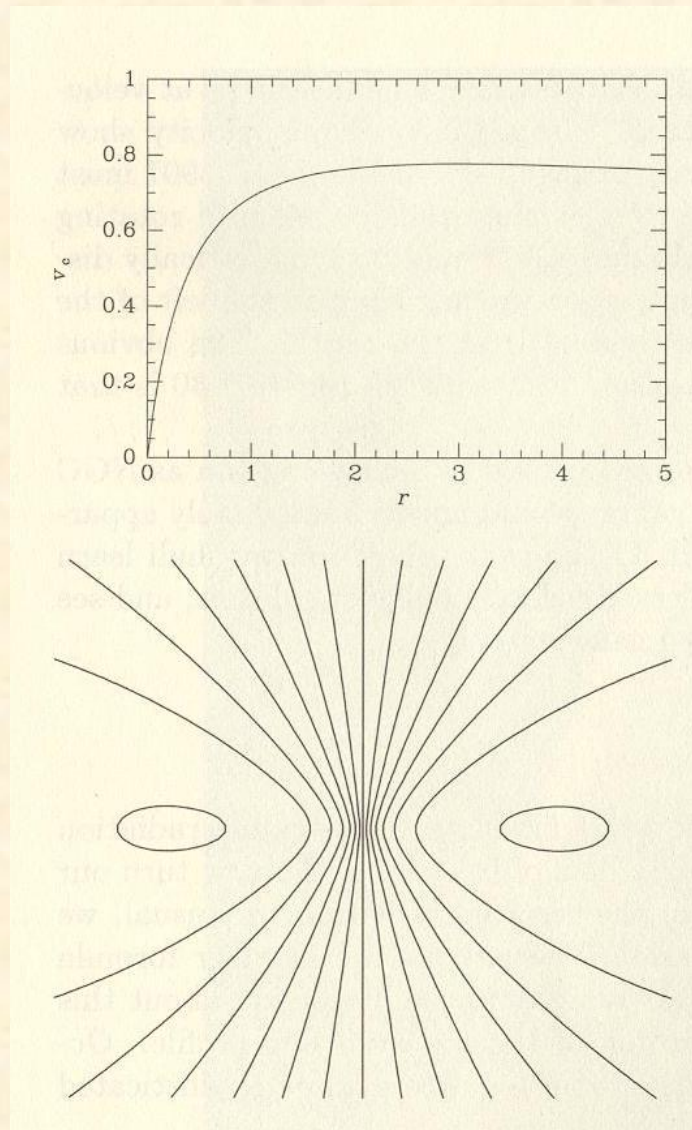
Velocity law



Spider Diagrams for $i = 60^\circ$
(isovelocity contours in equal steps)



More realistic example

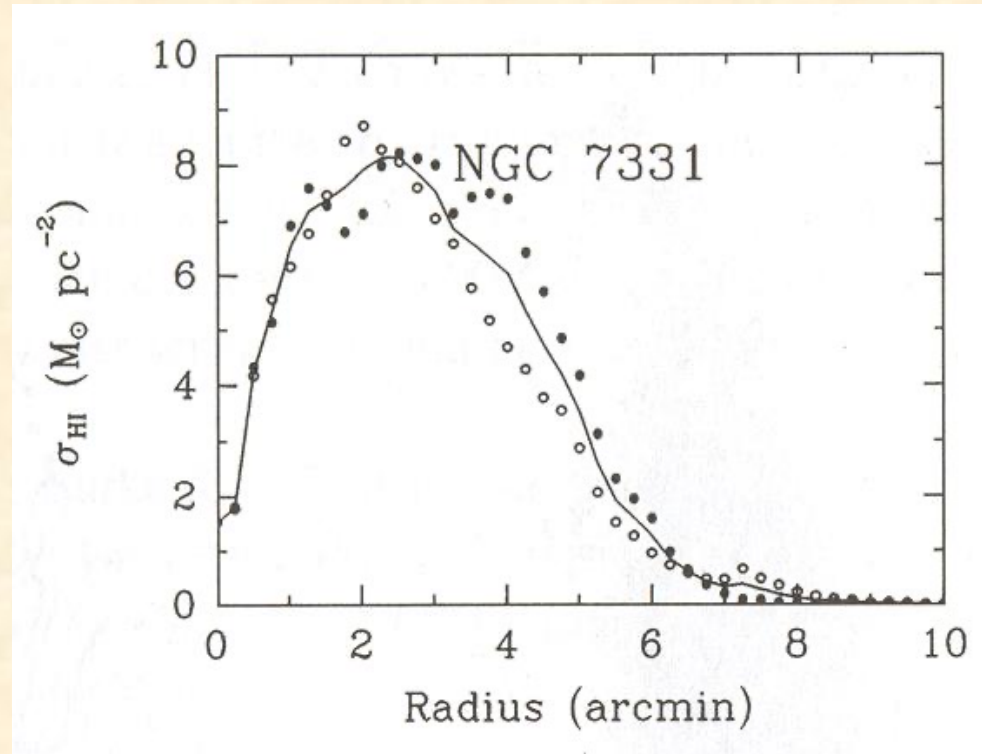


Binney & Merrifield, 506

Radial Velocities from H I 21-cm Emission

- H I gas is better tracer of kinematics than stars:
 - More uniformly distributed and more extensive

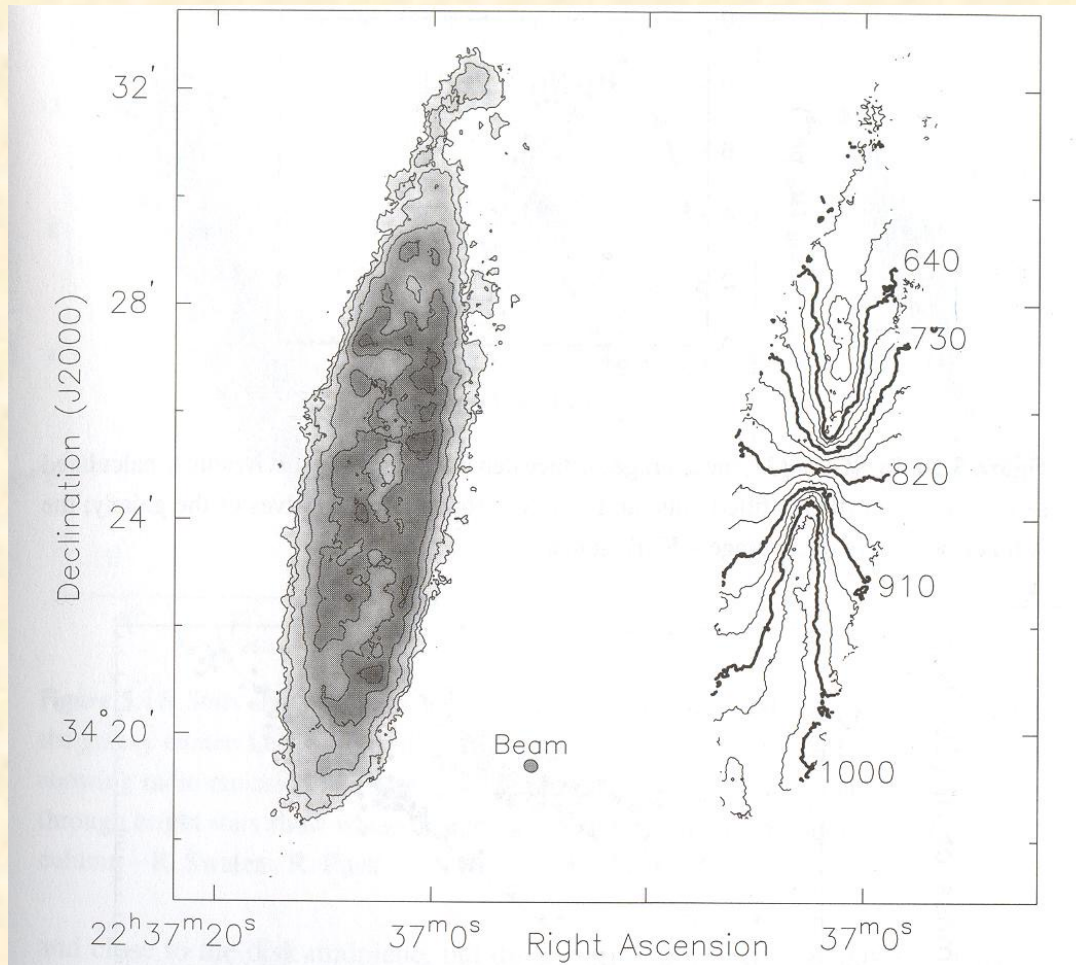
Surface Density



(Sparke & Gallagher,
p. 211)

- H I typically detected out to $2R_{25}$ to $4R_{25}$ (R_{25} of MW ≈ 10 kpc)
- Mass (H I) ≈ 1 to 10% Mass (stellar disk)
(Sa \rightarrow Sd)

NGC 7331 - HI Intensity and Spider Diagram

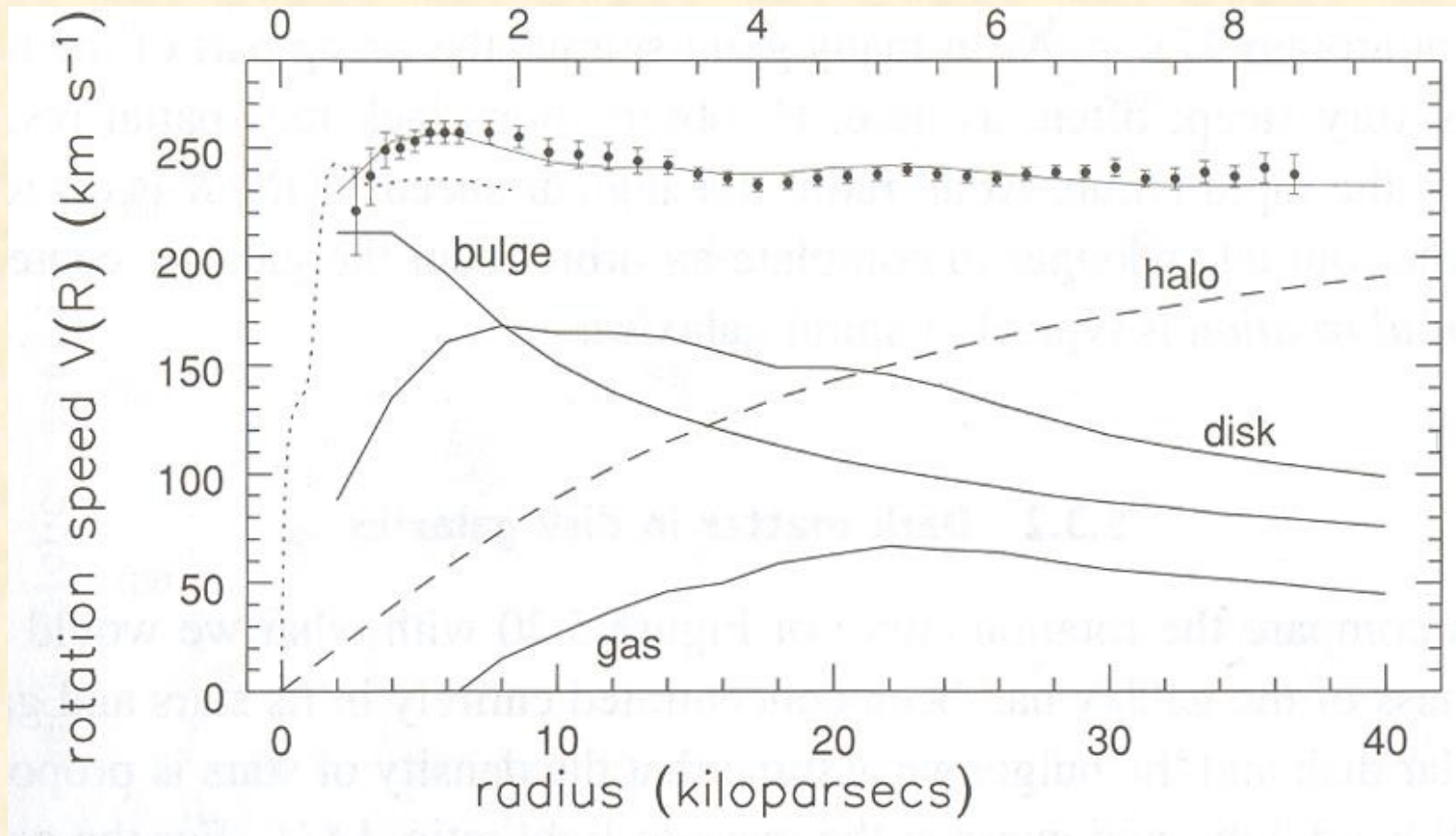


linear increase in v_r
followed by constant v_r

(Sparke & Gallagher, p. 210)

Complications: Isophotes twisted in same direction: warped disk
Gradient along minor axis – radial motion
Kinks in isophotes – random motions

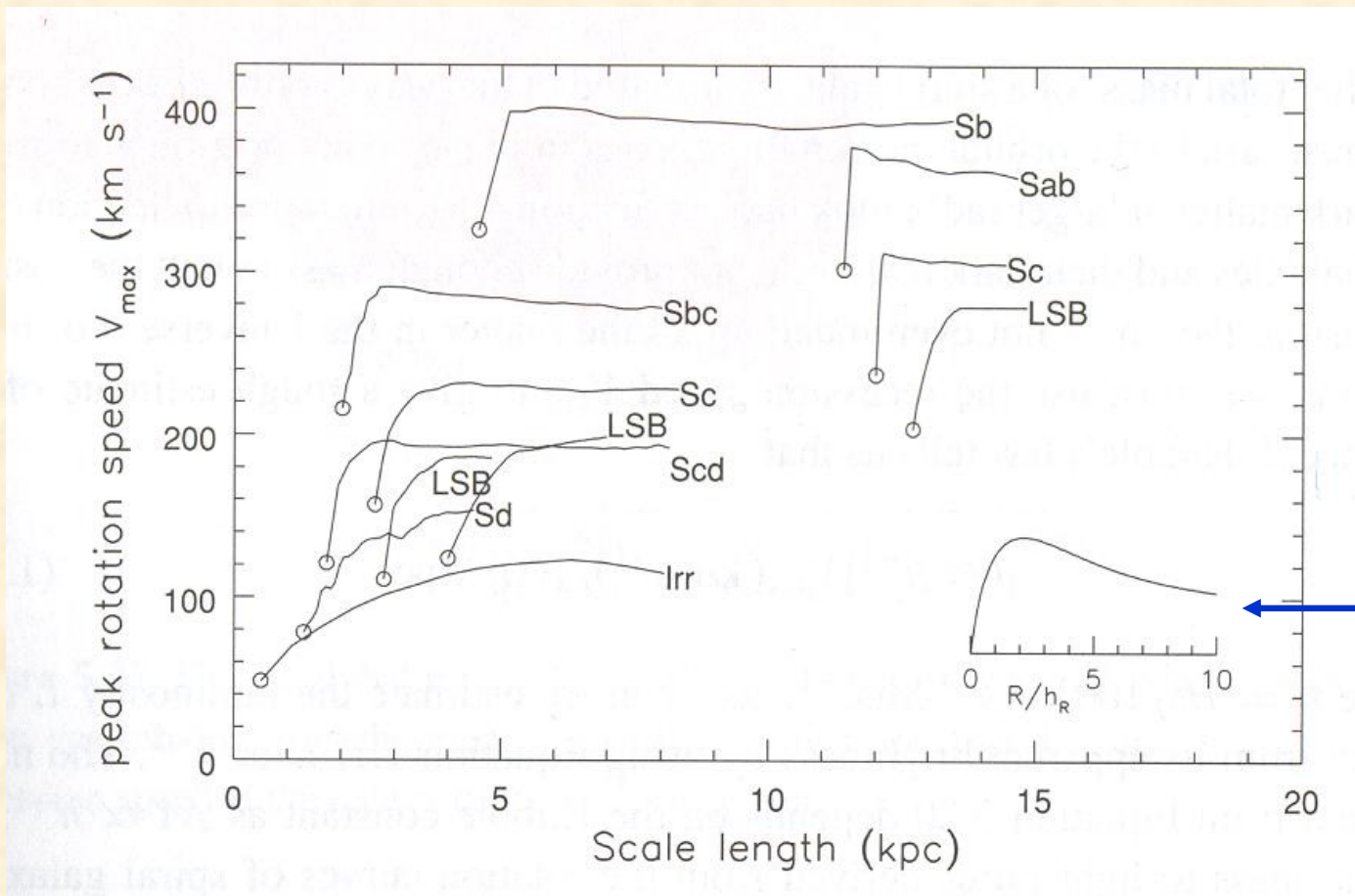
Rotation Curve (along major axis) for NGC 7331



(Sparke & Gallagher, p. 197)

- Dotted line: CO observations (traces colder molecular gas)
- Points and solid line: H I 21-cm measurements
- Bulge, disk, and gas: deduced from surface-brightness profiles
- Inferred dark halo mass: 2 to 4 times visible mass (in general)

Rotation Curves for Other Galaxies

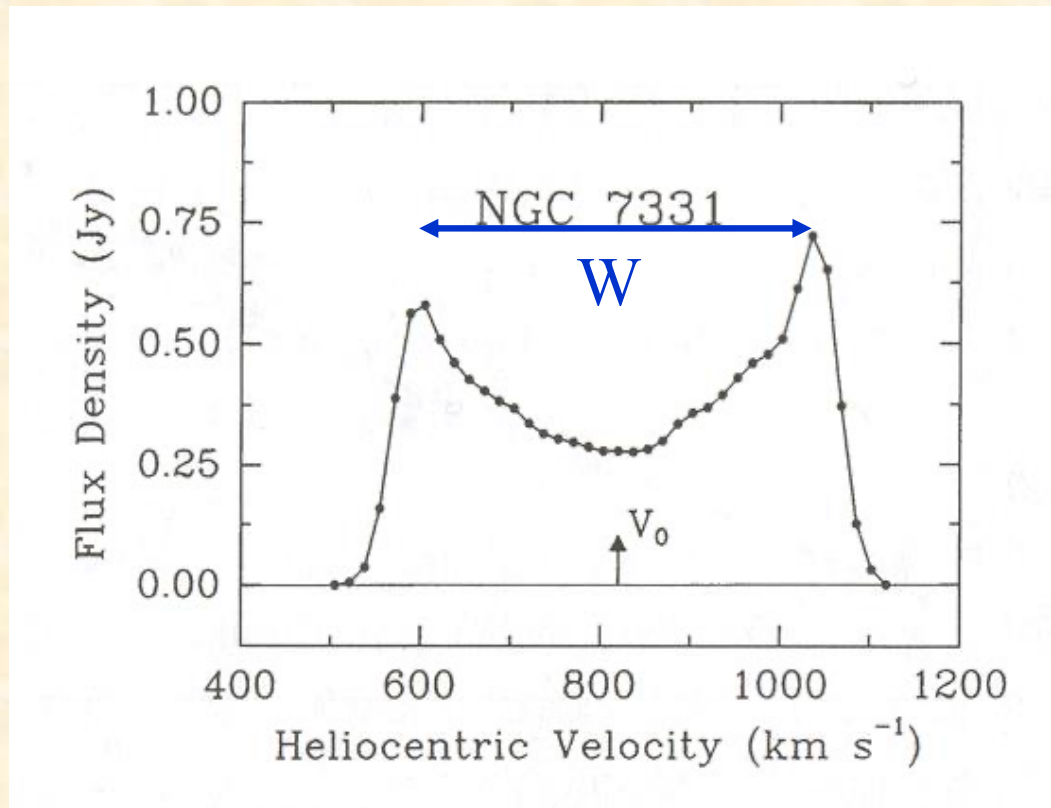


(Sparke & Gallagher, p. 218)

- Larger disk galaxies rotate faster
- Early types tend to rise more steeply
- Flat rotation curves: evidence for dark halos in disk galaxies

Tully-Fisher Relation

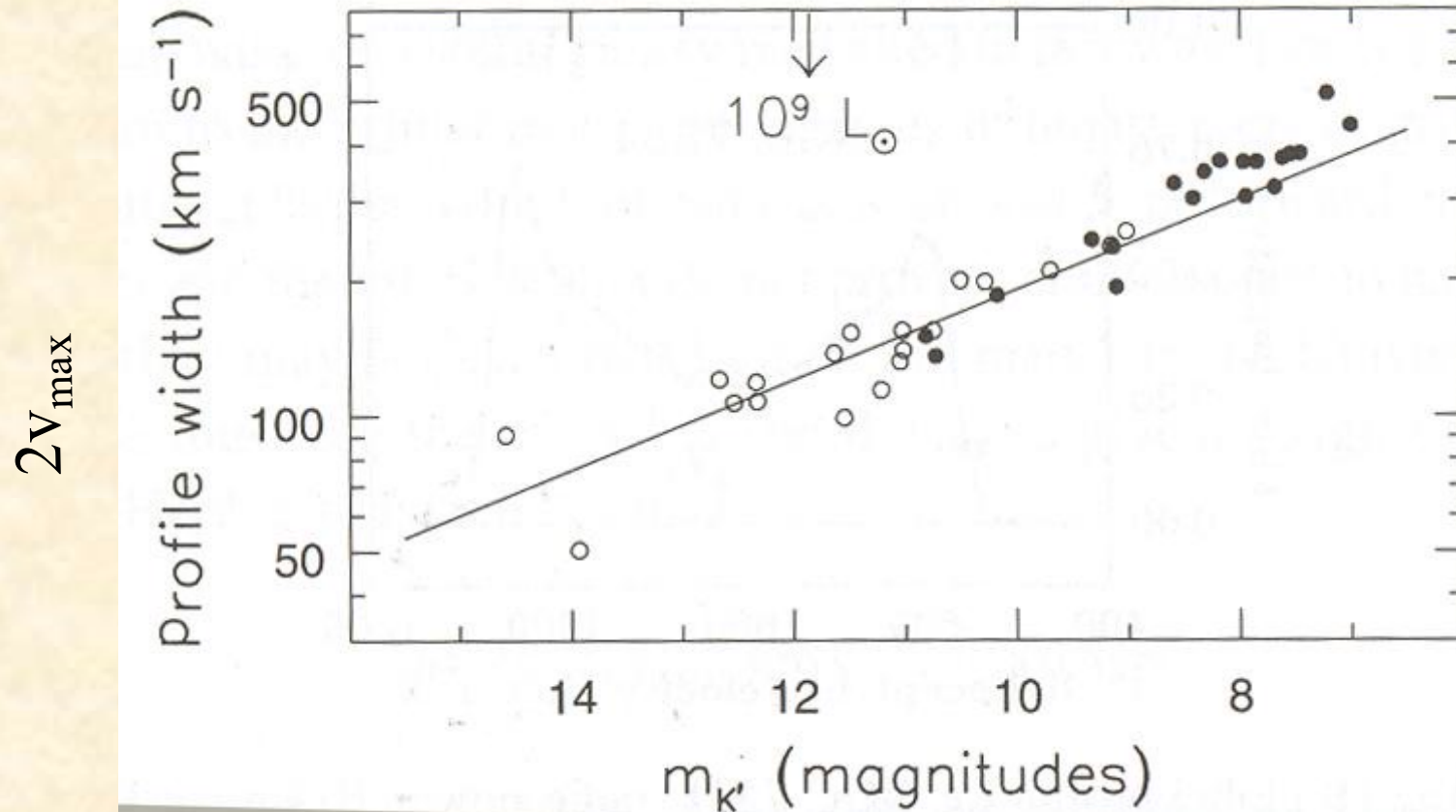
- Rotation curves not possible for more distant spirals
- Use “integrated” H I profile: double-horned common



(Sparke & Gallagher,
p. 220)

$$V_{\max} = \frac{1}{2}W/\sin(i)$$

Ex) Ursa Major Group



(Sparke & Gallagher, p. 221)

$$\text{Tully - Fisher : } L_{\text{IR}} \propto v_{\max}^4$$

$$\text{Recent calibration: } \frac{L_{\text{I}}}{4 \times 10^{10} L_{\text{I},\odot}} = \left(\frac{v_{\max}}{200 \text{ km s}^{-1}} \right)^4$$

Hand-waving Justification for Tully-Fisher:

$$M(R_d) \propto R_d V_{\max}^2$$

If M/L ratio is constant:

$$L \propto R_d V_{\max}^2$$

$$\text{Also : } L \propto I_0 R_d^2 \propto I_0 \frac{L^2}{V_{\max}^4}$$

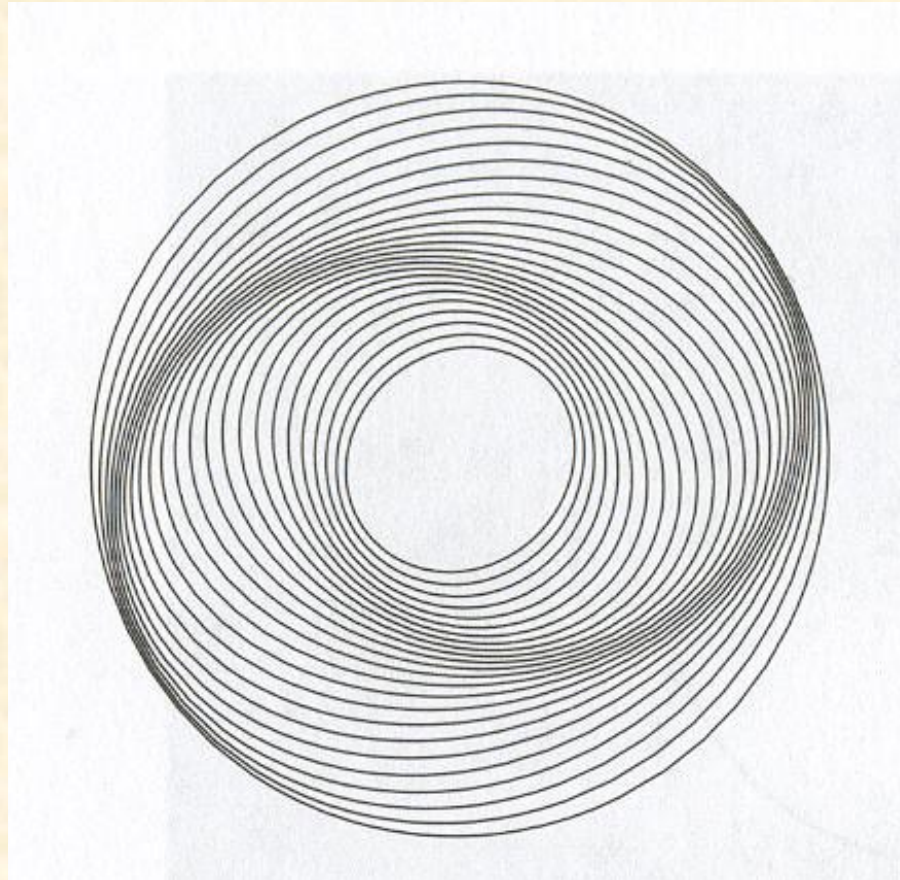
$$\text{If } I_0 \text{ is constant: } L \propto V_{\max}^4$$

This probably shouldn't work:

- I_0 and M/L are not constant with type or luminosity
- Velocities are affected by dark halo, luminosity is not

Spirals

- Even for constant $v(R)$, the angular velocity (v/R) drops with increasing distance
 - differential rotation should wind spirals up
- Theories:
 - 1) Starburst is stretched out by differential rotation:
 - works for fragmentary (flocculent) arms
 - 2) Density wave
 - continuous (including grand design) arms
 - pattern speed tends to be much slower than rotation
 - pattern is maintained by self gravity

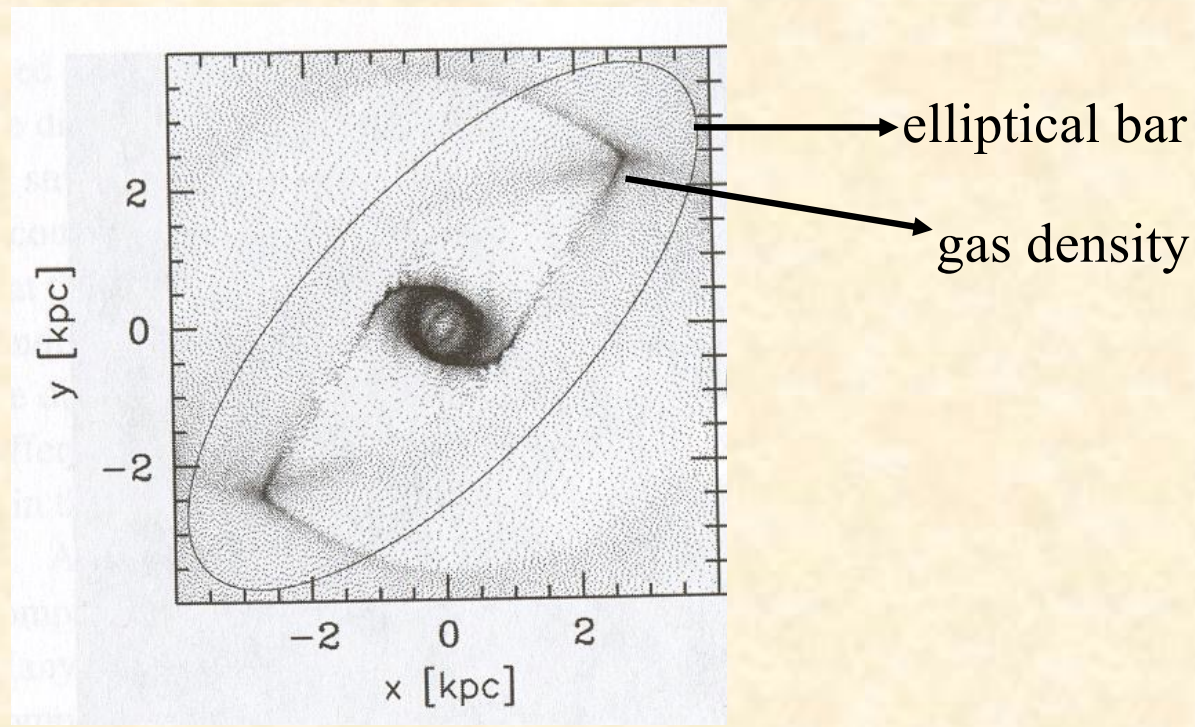


(Sparke & Gallagher, p. 219)

- Two-armed spiral: nested ovals with rotating position angles
- Originates from external (another galaxy) or internal (bar) perturbations

Bars

- Not a density wave - “stars remain in bars”
- As with spirals, pattern speed is slower than rotation (up to the co-rotation radius, where bar ends)
- Gas builds up and is shocked on leading edge → **infall**



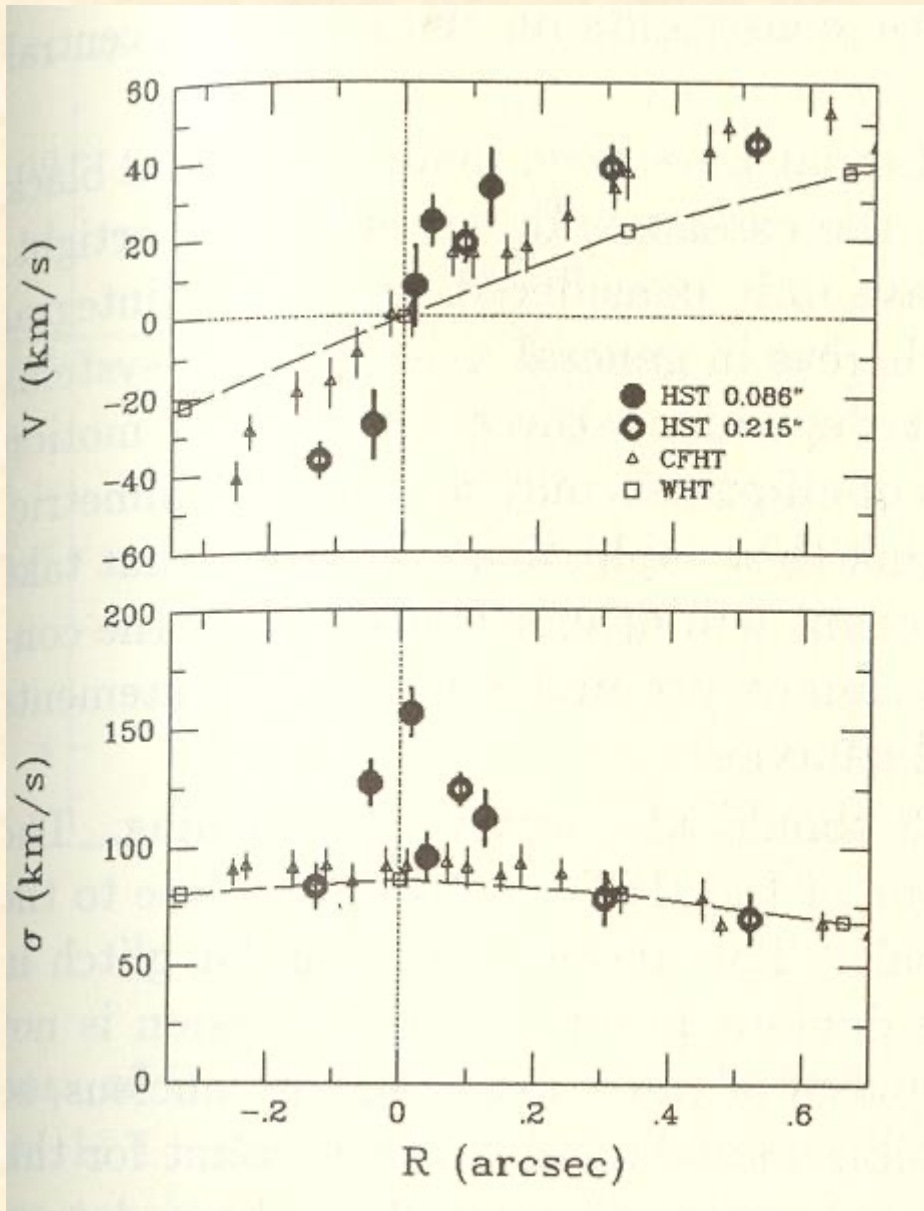
(Sparke & Gallagher, p. 235)

Supermassive Black Holes (SMBHs)

- SMBHs have been detected in the gravitational centers of most nearby galaxies.
- Direct methods to detect and measure masses of *quiescent* SMBHs are based on **resolved** spectroscopy:
 - 1) Stellar kinematics:
 - Should show rapid rise in v_r and/or σ near center.
 - Need high angular resolution (HST or Ground-based AO)
 - Use dynamical models (dominated by ellipticals and bulges).
 - What range of stellar orbits and mass density distributions $\rho(\mathbf{r})$ give the observed v_r , σ , and μ (2D) distributions?
 - Add a point-source mass (if necessary) to match the core.
 - 2) Measurement of positions, proper motions and radial velocities of individual stars
 - only the Milky Way $\rightarrow M_{\bullet} = 4.3 \times 10^6 M_{\odot}$

1) Stellar Kinematics from HST

Ex) M32 (compact dE)



- HST detected high v_r and σ in core.
- trends smeared out in ground-based telescopes
- SMBH Mass: $M_{\bullet} = 3 \times 10^6 M_{\odot}$

“Clincher”: STIS LOSVD in core shows high-velocity wings (Joseph, et al. 2001, ApJ, 550, 668)

Why do we need high angular resolution?

- What is the radius of influence for the SMBH in M32?

$$r = \frac{GM_{\bullet}}{\sigma_*^2}, \text{ where } \sigma_* = \text{typical stellar velocity dispersion}$$

For M32, $r \approx 1 \text{ pc} \rightarrow 0.3''$ at a distance of 725 kpc.

- SMBH “machine”: HST’s Space Telescope Imaging Spectrograph (STIS) – long slit, high resolution spectra
 - angular resolution $\sim 0.1''$, velocity resolution $\sim 30 \text{ km/s}$
 - measured SMBH masses in many nearby galaxies
- Note these observations do not **prove** existence of SMBHs
 - Ex) M32 mass concentrated within $\sim 0.3 \text{ pc}$:
 - $\rightarrow \rho \approx 10^{-15} \text{ g cm}^{-3}$! (pretty good vacuum)
 - \rightarrow previously: must rely on arguments that stars inside this volume would collide and eventually form a SMBH

What would prove the existence of a SMBH?

- Gravitationally redshifted emission from gas within a few times the Schwarzschild radius (R_s)

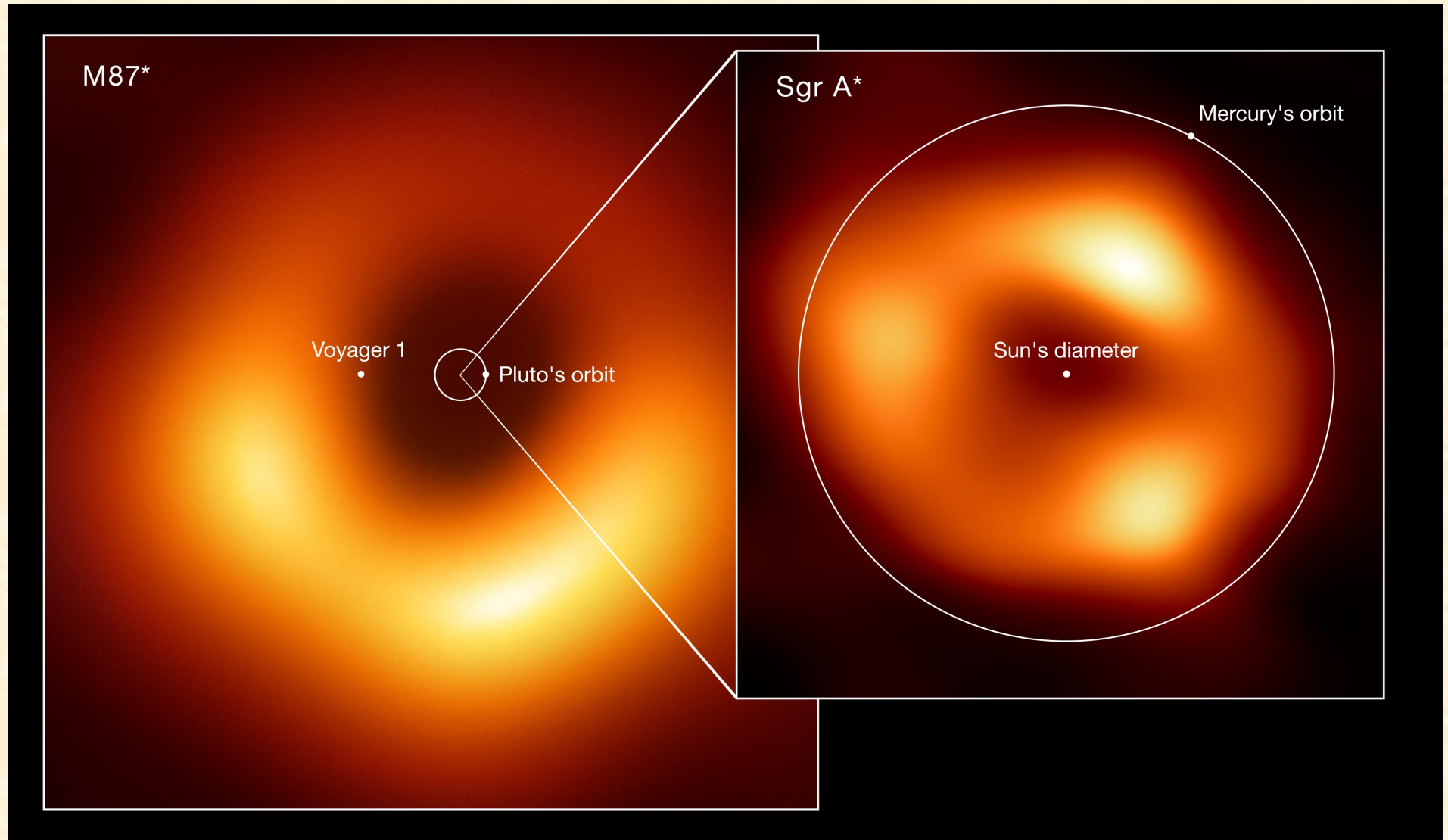
$$v_{\text{esc}}^2 = c^2 = \frac{2GM_{\bullet}}{R_s} \rightarrow R_s = \frac{2GM_{\bullet}}{c^2}$$

For M32: $R_s = 9 \times 10^{11} \text{cm} \approx 13R_{\odot}$

→ projected angular size: $\theta \approx 10^{-7} \text{arcsec}$

- No hope of resolving directly (can't get rotation curve)
- X-ray observations of AGN have detected **gravitationally-redshifted** Fe $K\alpha$ emission (presumably from accretion disk)
- Now we have proof from Event Horizon Telescope observations of M87 and Sgr A*.

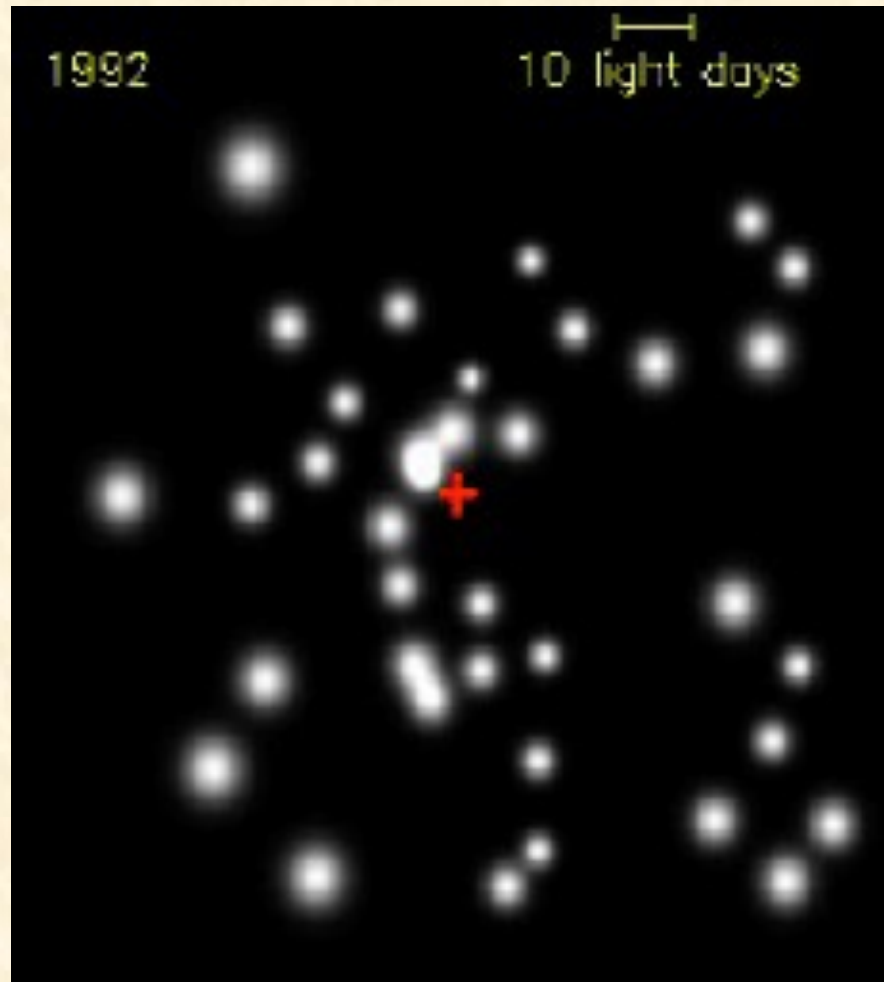
EHT



$$M_{\bullet} = 6.5 \times 10^9 M_{\odot}$$

$$M_{\bullet} = 4.3 \times 10^6 M_{\odot}$$

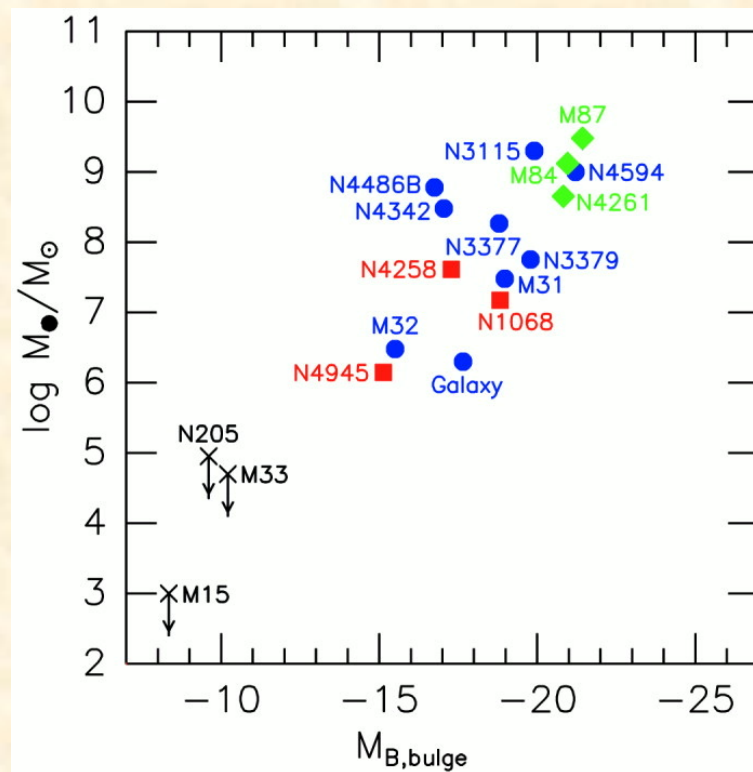
2) Individual Stars - Milky Way



- K band observations with NTT, VLT (mostly O and B supergiants)
- Proper motions plus radial velocities give $M_{\bullet} = 4.3 \times 10^6 M_{\odot}$

SMBH Mass/Bulge Correlations

- Kormendy et al. found a correlation between SMBH mass (M_{\bullet}) and absolute blue magnitude of the bulge/elliptical

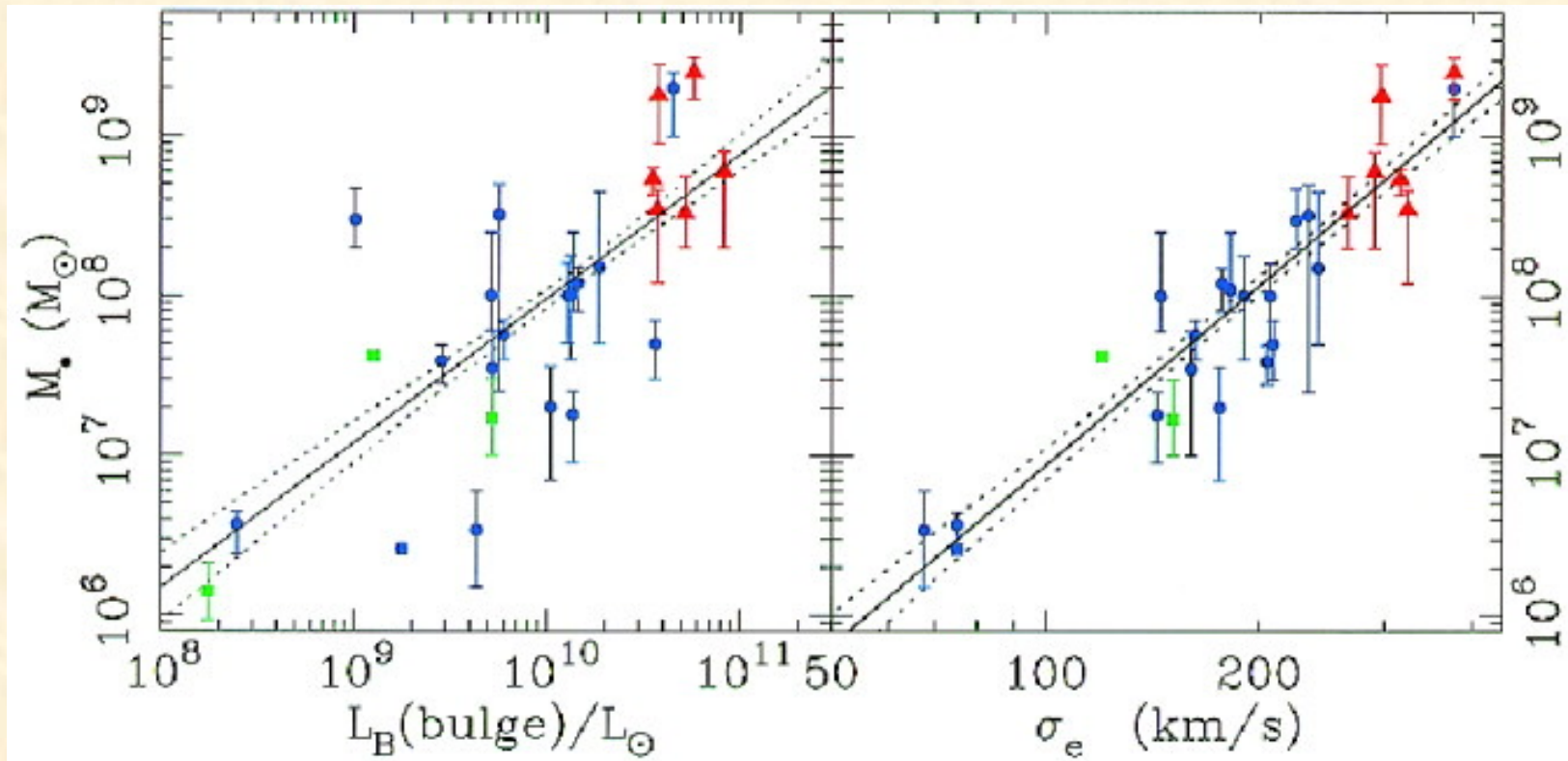


green – gas kinematics
blue – stellar kinematics
red – H₂O maser

(Kormendy, et al. 1998, AJ, 115, 1823)

- recent studies confirm: $L_{bulge} \sim M_{\bullet}$
- given a constant M/L ratio: $M_{\bullet} \approx 0.002 M_{bulge}$

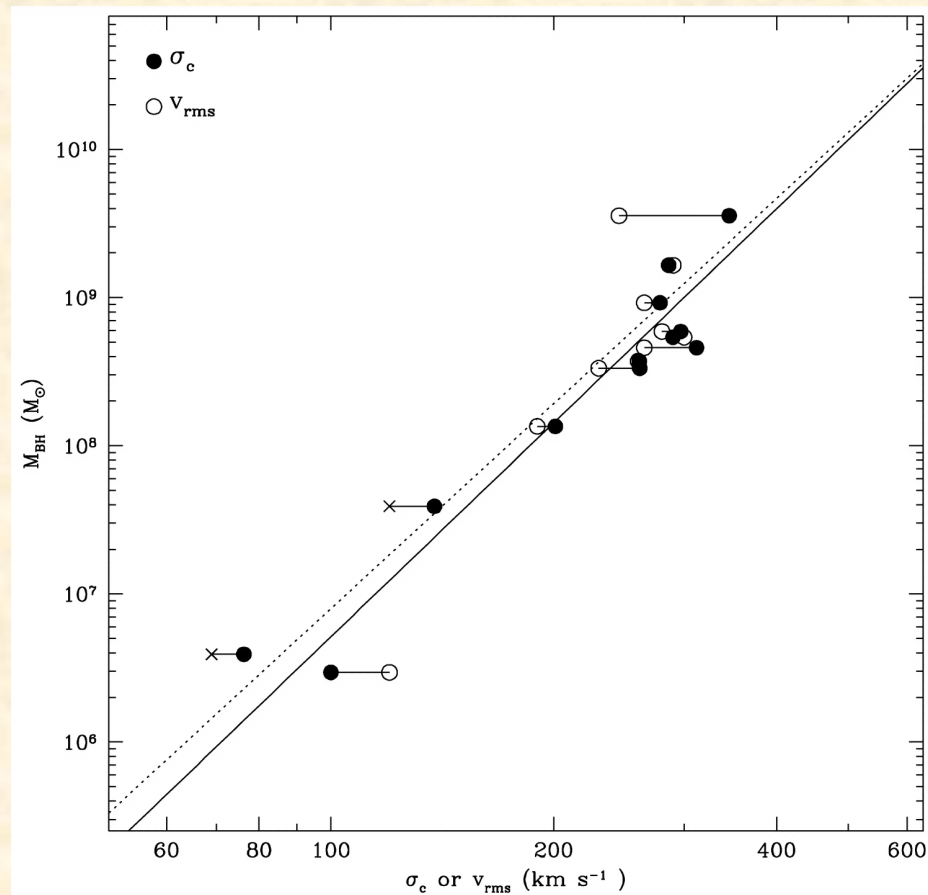
- Tighter correlations have been found with σ (bulge):
 - 1) $M_{\bullet} \sim \sigma^{3.75}$



(Gebhardt et al. 2000, ApJ, 539, L13)

- from stellar, gas kinematics, and masers
- σ_e : velocity dispersion of bulge within half-light radius

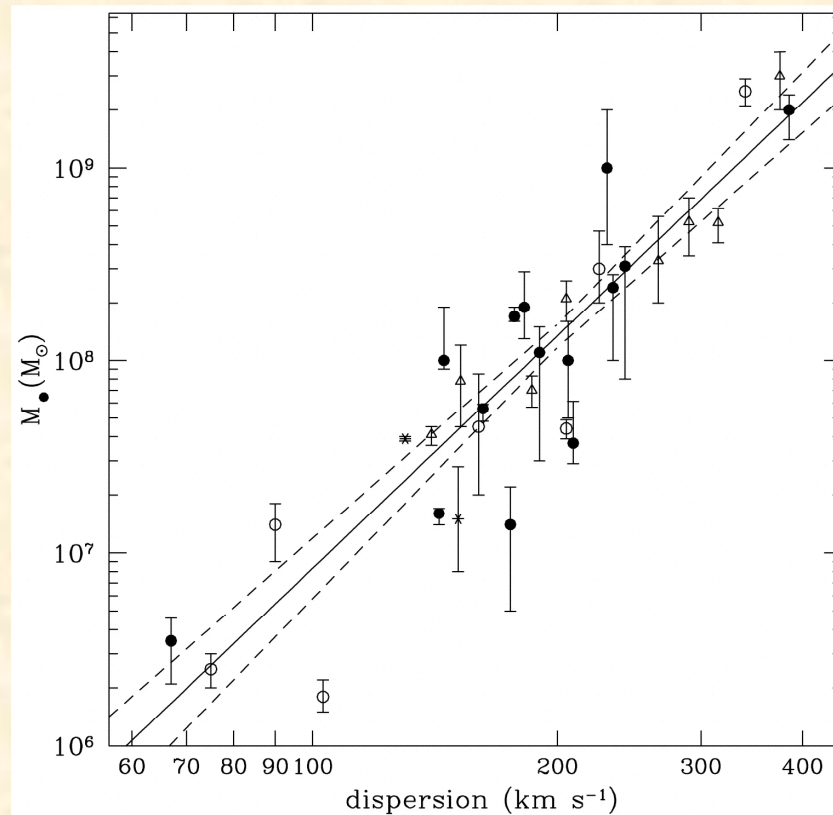
$$2) M_{\bullet} \sim \sigma^{4.80}$$



(Ferrarese & Merritt 2000, ApJ, 539, L9)

σ_c – velocity dispersion within 1/8
the effective radius of bulge

3) Subsequent Calibration: $M_{\bullet} \sim \sigma^4$



(Tremaine, et al. 2002, ApJ, 574, 740)

- based on stellar (circles), gas kinematics (triangles), masers (asterisks)
- previous disagreements probably due to different ways to measure $\sigma(\text{bulge})$

$$\log\left(\frac{M_{\bullet}}{M_{\oplus}}\right) = (4.02 \pm 0.32) \log\left(\frac{\sigma}{200 \text{ km s}^{-1}}\right) + (8.19 \pm 0.06)$$

Implications

- SMBHs present in all galaxies with a spheroidal component
 - For distant galaxies, M_{\bullet} can be inferred from σ_0 or L_{bulge}
 - SMBHs in AGN have the same mass as quiescent SMBHs for a given spheroidal (bulge) mass ($M_{\bullet} \approx 0.002 M_{\text{bulge}}$)
 - AGN were much more common in the past. Many quiescent SMBHs are dead remnants of AGN/QSOs.
 - How do SMBHs know about their bulges? – linked by evolution?
 - SMBHs formed by
 - 1) Overdensities in the early Universe?
 - 2) Massive Pop III stars?
 - 3) BHs from evolution of nuclear star cluster?
- Perhaps linked by AGN feedback: radiation and mass outflows determine size of SMBH and bulge