

Note on Debye length and plasma parameter

Debye Length

Introduction

The term plasma was introduced by Tonks and Langmuir (1929). The term 'plasma' is used to describe a region of an arc discharge where ion and electron densities are high and nearly equal. They compare the behaviour of discharge to blood, with ions acting like a rigid jelly when electrons oscillate. Plasma now refer to assemblies of charged particles studied in plasma physics, where matter both behaves as particle and fluid. This is similar to where matter behaves both as particles how quantum mechanics reveals dual particle-wave properties. The discussion adopts the gaussian unit system for electromagnetism, where electric and magnetic quantities are measured in electrostatic and electromagnetic units respectively.

From Maxwell's eqnⁿ

$$\nabla \times E = -\frac{1}{c} \frac{\partial B}{\partial t} \quad - \textcircled{1}$$

$$\nabla \cdot \mathbf{B} = 0 \quad - \quad (2)$$

$$\nabla \cdot \mathbf{E} = 4\pi \rho \quad - \quad (3)$$

$$\nabla \times \mathbf{B} = \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} + 4\pi \mathbf{j} \quad - \quad (4)$$

Where ρ is charge density, \mathbf{j} is current density, \mathbf{E} is electric field and \mathbf{B} is magnetic field.

So, from eqⁿ (1) and (2)

$$\frac{1}{c} \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{j} = 0 \quad - \quad (5)$$

Let us consider simple plasma composed of fully ionized hydrogen. Then mass density, charge density and current density is given by

$$\text{Mass density } (\rho) = n_e m_e + n_i m_i$$

$$\text{charge density } (\rho) = e(n_i - n_e)$$

$$\text{Current density } (\mathbf{j}) = \frac{e}{c} (n_i \mathbf{v}_i - n_e \mathbf{v}_e)$$

Where, n_e, n_i = no. of density of electron and Proton
 m_e, m_i = mass of electron and Proton
 $-e, e$ = charge of electron and Proton.

Let us consider a plasma with initially uniform proton and electron densities, with no net charge density, so initially electric field is zero. Let n_0 be initial densities of Proton and electron. Let us suppose proton density change from n_0 to $(1-\delta)n_0$ in region $-L < x < L$. If L is sufficiently small

the electric field due to this charge will be small
 So, electron will suffer negligible distribution. If
 L is very large then distribution of electron will
 be very huge. If we take value of L so that such
 a way that just transition takes place then eqnⁿ
 (3) becomes

$$\frac{d^2\phi}{dn^2} = 4\pi\delta n_0 e - \quad (6)$$

Consider that plasma as a whole is maintained
 at potential $\phi=0$, then eqnⁿ (6) becomes

$$\left. \begin{aligned} \phi &= 2\pi\delta n_0 e (\lambda^2 - L^2) & |x| < L \\ \phi &= 0 & |x| > L \end{aligned} \right\} - \quad (7)$$

So,

$$\phi(0) = -2\pi\delta n_0 e L^2 - \quad (8)$$

Consider, plasma has Temp^r T_0 . K.E. is $\frac{1}{2}kT$ in
 each degree of freedom. If $\phi(0)$ is so small that
 avg. thermal energy ≈ 0 , then there will be only
 small change in state of plasma. If $\phi(0)$ is very
 large then only small few electron can reach $x=0$

So, for 'quasi neutral' state

$$\frac{1}{2}kT > 2\pi\delta n_0 e^2 L^2 - \quad (9)$$

Also,

$$\delta < (\lambda_D/L)^2, \quad \lambda_D = \frac{kT}{4\pi n e^2} - \quad (10)$$

Also, For Plasma Parameter (Λ)

$$\lambda_D = \sqrt{\frac{kT}{4\pi n e^2}} \quad \text{--- (ii) OR } \lambda_D = \sqrt{\frac{\epsilon kT}{n e^2}}$$

$$\lambda_D = 10^{0.84} n^{-1/2} T^{1/2}$$

Where,

λ_D is Debye length

It's

(1) $\delta = -\lambda_D$ proton is completely removed from
- $L < n < L$, plasma is quasi neutral
($L \ll \lambda_D$)

(2) If $L \gg \lambda_D \rightarrow$ behaves as vacuum.

Height Scale

Let us consider solar corona, in which $m_p \gg m_e$
gravitational field only acts on proton only. Due
to gravitational field densities decreases exponentially
with height, which is given by

$$H = \frac{kT}{m_{av} \cdot g}$$

m_{av} avg. Particle mass $\approx \frac{1}{2} m_p$
 g is gravi-acceleration.

Ref:- Plasma phy by: Peter A. Sturrock
- Arhab R. Choudhori

Also, For Plasma Parameter (Λ)

$$\Lambda = \frac{4\pi}{3} \lambda_D^3 n_e$$

Where, $\Lambda_e = n_e \lambda_D^3$, $\Lambda_s = n_s \lambda_{Ds}^3$

i.e. $\Lambda_e = \left(\frac{k}{4\pi e^2} \right)^{3/2} n_e^{-1/2} T_e^{3/2}$

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How - 3

Calculate the numerical value of the Debye length and Plasma Parameter in

- (a) the solar corona ($T = 10^6 \text{K}$, $n_e \sim 10^{14} \text{m}^{-3}$)
- (b) the earth magnetosphere ($T = 10^7 \text{K}$, $n_e \sim 10^7 \text{m}^{-3}$)
- (c) the interior of the sun ($T \sim 10^3 \text{K}$, $n_e \sim 10^7 \text{m}^{-3}$)
- (d) the interior of white dwarf (10^7K , $n_e \sim 29 \times 10^{32} \text{m}^{-3}$)
- (e) the interior of a neutron star. (10^{11}K to 10^{13}K , $n_e \sim 10^{40} \text{m}^{-3}$)

Soluⁿ

$$\lambda_D (\text{solar corona}) = \sqrt{\frac{\epsilon K T}{n e^2}}$$

Where

$$\epsilon = 8.85 \times 10^{-12} \text{Fm}^{-1}$$

$$K_B = 1.38 \times 10^{-23} \text{J K}^{-1}$$

$$\lambda_D = \sqrt{\frac{8.85 \times 10^{-12} \times 10^6 \times 1.38 \times 10^{-23}}{10^{14} \times (1.67 \times 10^{-19})^2}}$$

$$= \sqrt{4.07 \times 10^{-6} \text{m}} = 2.18 \times 10^{-3} \text{m} = 0.218 \text{cm}$$

$$\lambda_D = 2.18 \text{m} = 0.218 \text{cm}.$$

$$\therefore \lambda_D = 0.218 \text{cm}$$

$$\text{Plasma Parameter } (N) = \frac{4\pi}{3} \lambda_D^3 \cdot n_e$$

$$= \frac{4\pi}{3} \times (2.18 \times 10^{-3})^3 \times 10^{15}$$

$$= 4.3 \times 10^7$$

(b) Earth magnetosphere ($T \sim 10^7 \text{ K}$, $n_e \sim 10^7 \text{ m}^{-3}$)

$$\lambda_D = \sqrt{\frac{\epsilon K_B T}{n_e e^2}}$$

$$= \sqrt{\frac{8.85 \times 10^{-12} \times 1.38 \times 10^{-23} \times 10^7}{10^7 \times (1.69 \times 10^{-19})^2}}$$

$$= \sqrt{\frac{12.0213 \times 10^3}{2.8561}}$$

$$= \sqrt[3]{4.2 \times 10^3 \text{ m}}$$

$$= 64.80 \text{ m}$$

$$\text{Plasma Parameter } (N) = \frac{4\pi}{3} \lambda_D^3 \cdot n_e$$

$$= \frac{4\pi}{3} \cdot (64.80)^3 \cdot 10^7$$

$$= 1.13 \times 10^{13}$$

(c) the interior of the sun ($T \sim 10^7 \text{ K}$, $n_e \sim 10^{30} \text{ m}^{-3}$)

$$\lambda_D = \sqrt{\frac{8.85 \times 10^{-12} \times 10^7 \times 1.38 \times 10^{-23}}{10^{30} \times (1.69 \times 10^{-19})^2}}$$

$$= 6.91 \times 10^{-10} \text{ m}$$

$$\text{Plasma Parameter}(\Lambda) = \frac{4\pi}{3} \lambda_D^3 n_e$$

$$= \frac{4\pi}{3} \times (6.91 \times 10^{-10})^3 \times 10^{30}$$

$$= 1.38 \times 10^{20}$$

(d) The interior of white dwarf

$$\lambda_D = \sqrt{\frac{8.85 \times 10^{-12} \times 10^7 \times 1.38 \times 10^{-23}}{10^{32} \times (1.69 \times 10^{-19})^2}}$$

$$\left. \begin{array}{l} T = 10^7 \text{ K} \\ n_e = 10^{32} \end{array} \right\}$$

$$= \sqrt{\frac{8.85 \times 1.38 \times 10^{-12+7-23}}{(1.69)^2 \times 10^{-38+32}}}$$

$$= \sqrt{\frac{1.22 \times 10^{-27}}{2.56 \times 10^{-6}}} = 6.9 \times 10^{-11} \text{ m}$$

$$\text{Plasma Parameter}(\Lambda) = \frac{4\pi}{3} (\lambda_D)^3 n_e$$

$$= 1.38 \times 10^2$$

H.W # 4

We have

Momentum eqn for electron

$$m_e n \frac{d\vec{v}_e}{dt} = nq \left(\frac{\vec{v}_e \times \vec{B}}{c} \right) + nqE - \nabla q_e \quad (1)$$

Momentum eqn for Proton

$$m_p n \frac{d\vec{v}_p}{dt} = -nq \left(\frac{\vec{v}_p \times \vec{B}}{c} \right) - nqE - \nabla q_p \quad (2)$$

We know that,

Total pressure of gas is equal to ^{sum of} pressure exerted by electron and Proton

$$\text{i.e. } P = P_e + P_p \quad (3)$$

from the definition of current density

$$(\vec{J}) = qn_e \vec{v}_e - qn_p \vec{v}_p \quad (4)$$

now, Add eqn (1) and (2)

$$m \left(m_e \frac{d\vec{v}_e}{dt} + m_p \frac{d\vec{v}_p}{dt} \right) = nq \left(\frac{\vec{v}_e \times \vec{B}}{c} - \frac{\vec{v}_p \times \vec{B}}{c} \right) \quad (5)$$

Since, Pressure exerted by electron is very small as compare to pressure exerted by Proton so from eqn (5) we neglect mass of electron term.

$$m_p \frac{d\vec{v}_p}{dt} = \vec{J} \times \frac{\vec{B}}{c} - \vec{\nabla} p \quad (6)$$

Also, we have relation

$$\vec{v} = \frac{n e m_e v_e + m_p m_p v_p}{m_e + m_p}$$

dividing by m_e

$$n v = n_p v_p \quad (7)$$

So, eqⁿ (6) can be written as

$$\boxed{\frac{S d\vec{v}}{d\epsilon} = \vec{J} \times \frac{\vec{B}}{c} - \vec{\nabla} p} \quad (8)$$

from subtracting eqⁿ (1) - (11), we get

$$\boxed{\vec{\nabla} \times \frac{\vec{B}}{c} + \vec{E} = 0} \quad (9)$$

by using eqⁿ (9)

$$E = -(\vec{\nabla} \times \vec{B}) \quad (10)$$

Substituting (10) in (2)

$$m_n \frac{d\vec{v}}{dt} = -nq \left[\frac{\vec{v} \times \vec{B}}{c} \right] + nqE - \vec{\nabla} p$$

$$m_n \frac{d\vec{v}}{dt} = -nq \left[\frac{\vec{v} \times \vec{B}}{c} \right] + nq \left[\frac{\vec{v} \times \vec{B}}{c} \right] - \vec{\nabla} p$$

$$\boxed{m_n \frac{d\vec{v}}{dt} = -\vec{\nabla} p}$$

(11) - momentum eqⁿ.

Mass conservation eqnⁿ

$$\boxed{\frac{ds}{dt} + \nabla \cdot (s\mathbf{v}) = 0} \quad - (12)$$

③ Induction eqnⁿ

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) \quad (\text{from 10})$$

$$\boxed{\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}} \quad - (13)$$

So, MHD eqnⁿ with eliminating electric field are of eqnⁿ (11) (12) and (13)